

KAM for beating solutions of the quintic NLS

E. Haus*, M. Procesi**

* Dipartimento di Matematica e Applicazioni “R. Caccioppoli”
Università degli Studi di Napoli “Federico II”, Napoli, Italy

** Dipartimento di Matematica e Fisica
Università di Roma Tre, Roma, Italy

E-mail: emanuele.haus@unina.it, procesi@mat.uniroma3.it

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Abstract

We consider the nonlinear Schrödinger equation of degree five on the circle $\mathbb{T} = \mathbb{R}/2\pi$. We prove the existence of quasi-periodic solutions with four frequencies which bifurcate from “resonant” solutions (studied in [18]) of the system obtained by truncating the Hamiltonian after one step of Birkhoff normal form, exhibiting recurrent exchange of energy between some Fourier modes. The existence of these quasi-periodic solutions is a purely nonlinear effect.

1 Introduction

We consider the quintic NLS on the circle

$$-i\partial_t u + \partial_{xx} u = \sigma |u|^4 u, \quad u = u(t, x) \quad (t, x) \in \mathbb{R} \times \mathbb{T} \quad \mathbb{T} := \mathbb{R}/(2\pi) \quad (1.1)$$

where $\sigma = \pm 1$. This is an infinite dimensional dynamical system with Hamiltonian

$$\mathcal{H} = \int_{\mathbb{T}} |\partial_x u|^2 + \frac{\sigma}{3} \int_{\mathbb{T}} |u|^6 \quad (1.2)$$

having the mass (the L^2 norm) and the momentum

$$\mathbb{L} = \int_{\mathbb{T}} |u|^2, \quad \mathbb{M} = \int_{\mathbb{T}} \text{Im}(u \partial_x u) \quad (1.3)$$

as constants of motion. Classical results imply that, for small values of the mass, equation (1.1) is globally well-posed in H^s with $s \geq 1$. This is a consequence of the local well-posedness result in [6], combined with the conservation law (1.2) (if the mass is small,

the Hamiltonian controls the H^1 -norm also for $\sigma = -1$, via the Gagliardo-Nirenberg inequality). However, one expects solutions of (1.1) to exhibit a rich variety of qualitative behaviors. Even close to zero, very little is known on the qualitative structure of typical solutions and the literature is confined to the study of special solutions which exhibit interesting features. A fruitful approach is to take the dynamical systems point of view and to apply the powerful tools of *singular perturbation theory*, such as KAM theory, Birkhoff Normal Form, Arnold diffusion, first developed in order to study finite-dimensional systems.

A first elementary feature is that NLS equation has an elliptic fixed point at $u = 0$. However such fixed point is completely resonant (i.e. all the linear solutions are periodic). It is well known, already in a finite dimensional setting, that such resonances may produce hyperbolic tori, secondary tori and energy transfer phenomena where the solutions of the non-linear system differ drastically from the ones of the linear system.

The dynamical systems approach was first applied in the early '90s to model PDEs on $[0, \pi]$ with Dirichlet boundary conditions in order to prove existence of quasi-periodic solutions and finite dimensional reducible KAM tori (see Definition 1.1 and 1.2). Among the vast literature we mention [33], [24], [25] and [7],[13], [10] for the case of periodic boundary conditions. We note that the above papers all take advantage of some simplifying assumptions and hence they do not apply to our setting. The first results on reducible KAM tori for NLS equations were proved in [15] for equations with external parameters, see also [31]. Reducible KAM tori for (1.1) were then proved in [26] and [16], see also [27],[29] and [32]. Differently from the (integrable) cubic case, it can be seen that for appropriate choices of the initial data such solutions have a finite number of linearly unstable directions. Then it is quite natural to wonder whether one may prove the existence of an orbit which starts close to the unstable manifold of one solution and in very long time arrives close to the stable manifold of a different solution, thus exhibiting a large drift in the action variables. This approach was first proposed in [11] in order to study the growth of Sobolev norms for the cubic NLS on \mathbb{T}^2 (see [20] for an estimate of the time) and then generalized to any analytic NLS on \mathbb{T}^2 in [22] and [19]. Unfortunately the approach proposed in these papers fails equation (1.1) on the circle, see Appendix C of [22] for a discussion of the problem. The main point is that even if it is still true that unstable solutions exist one is not able to prove full energy transfer. In [18] the authors study a quintic NLS on the circle as (1.1) and prove the existence of *beating* solutions which exhibit a *recurrent energy exchange* between some Fourier modes for a *long but finite* time interval (see also [23]). The result follows by showing that the Birkhoff normal form Hamiltonian exhibits such a behavior and that its dynamics approximates the true one for long times. It is then very natural to ask whether such solutions persist for infinite times, this is the question assessed in the present paper. The solutions that we construct are genuinely non-linear: they do not bifurcate from linear solutions, which do not feature any transfer of energy between Fourier modes. Regarding other genuinely non-linear periodic solutions for PDEs we mention the papers [2] and [5].

It is interesting to note that our quasi-periodic solutions correspond to *secondary* KAM tori (see for instance [4] for a discussion about secondary KAM tori in the finite-dimensional context), since they are 4-dimensional, but homotopically equivalent to 3-dimensional tori in the phase space. To the best of our knowledge, this is the first result establishing

existence of secondary KAM tori for PDEs. It would also be interesting to reproduce a similar result for the cubic NLS on \mathbb{T}^2 : we believe that this can be done, however the extension of this type of result to a higher dimensional setting involves some non-trivial technical difficulties.

Among other results proving long time existence of beating solutions, we mention the work [17], dealing with two coupled cubic Schrödinger equations on the circle. Moreover, the modified scattering argument proved in [21] implies that the cubic NLS on the product space $\mathbb{T}^2 \times \mathbb{R}$ has solutions exhibiting beatings over an infinite forward (or backward) time.

Main results and strategy of the paper. We start by defining the main objects on which our result relies.

Definition 1.1. *Given a rationally independent¹ $\omega \in \mathbb{R}^n$, a quasi-periodic solution of frequency ω is a torus embedding $\mathfrak{I} : \mathbb{T}^n \ni \varphi \mapsto u(\varphi, x)$ such that $u(\omega t, x)$ is a classical solution for (1.1).*

Definition 1.2 (reducible KAM torus). *A quasi-periodic solution of frequency ω is a reducible KAM torus if (in appropriate complex symplectic variables φ, y, z, \bar{z}) it is expressed by the equations $y = z = 0$ and moreover the associated Hamiltonian vector field X_H restricted to the torus is $\sum_{i=1}^n \omega(\xi) \partial_{\varphi_i}$, while X_H linearized at the torus is block-diagonal in the normal variables with x -independent block matrices of uniformly bounded dimension.*

From now on, we shall systematically work on the Hamiltonian (1.2), which we study in the Fourier coefficients $u(t, x) = \sum_{j \in \mathbb{Z}} u_j(t) e^{ijx}$ where the sequence $\{u_j\}_{j \in \mathbb{Z}} \in \ell^{a,p}$.

$$\mathcal{H} = \sum_{j \in \mathbb{Z}} j^2 |u_j|^2 + \sigma \sum_{j_1+j_2+j_3=j_4+j_5+j_6} u_{j_1} u_{j_2} u_{j_3} \bar{u}_{j_4} \bar{u}_{j_5} \bar{u}_{j_6} \quad (1.4)$$

with the constants of motion

$$\mathbb{L} = \sum_{j \in \mathbb{Z}} |u_j|^2, \quad \mathbb{M} = \sum_{j \in \mathbb{Z}} j |u_j|^2. \quad (1.5)$$

Remark 1.1. *For all $k \in \mathbb{Z}$, equation (1.1) is invariant with respect to the transformation*

$$v_j(t) := u_{j+k}(t) e^{-it(k^2+2jk)}. \quad (1.6)$$

Since the actions satisfy $|v_j|^2 = |u_{j+k}|^2$, this means that for any given solution $u(t, x)$ of (1.1), there exists another solution $v(t, x)$ whose actions have the same dynamics, up to a translation by k of the Fourier support.

Having fixed our notation, the main result of the paper is the following.

Theorem 1.1. *There exist $\varepsilon_0 > 0$ and a compact domain $\mathcal{O}_0 \subset \mathbb{R}^4$ such that, for all $0 < \varepsilon < \varepsilon_0$, there exists $\mathcal{O}_\varepsilon \subset \mathcal{O}_0$, satisfying $|\mathcal{O}_\varepsilon|/|\mathcal{O}_0| \rightarrow 1$ as $\varepsilon \rightarrow 0$, and a locally invertible Lipschitz mapping $\omega^\infty : \mathcal{O}_\varepsilon \rightarrow \mathbb{R}^4$, with the following property. Fix $k \in \mathbb{Z}$ and let*

¹i.e. a vector $\omega \in \mathbb{R}^n$ such that for all $k \in \mathbb{Z}^n \setminus \{0\}$ one has $\omega \cdot k \neq 0$.

$\mathcal{S} := \{-2, -1, 1, 2\}$. For all $\xi = (\xi_0, \xi_1, \xi_2, \xi_3) \in \mathcal{O}_\varepsilon$, equation (1.1) admits a quasi-periodic solution²

$$u_{\varepsilon,k}(\xi; t, x) = \varepsilon^{1/4} \left(\sum_{j \in \mathcal{S}} v_j(\xi; \omega^\infty(\xi)t) e^{i(j+k)x} e^{i(jk+k^2)t} + O(\varepsilon^{2/3^-}) \right) \quad (1.7)$$

of frequency $\omega^\infty(\xi) = (0, 1, 1, 4) + O(\varepsilon)$, such that the v_j satisfy

$$\begin{aligned} |v_j(\xi; \varphi_0, \varphi_1, \varphi_2, \varphi_3)|^2 &= \mathbf{f}_j(\xi; \varphi_0) \\ \mathbf{f}_{-2}(\xi; 0) &\leq \frac{1}{2} \quad \mathbf{f}_1(\xi; 0) \leq 1 \quad \mathbf{f}_{-1}(\xi; 0) \geq 3 \quad \mathbf{f}_2(\xi; 0) \geq \frac{3}{2} \\ \mathbf{f}_{-2}(\xi; \pi) &\geq \frac{3}{2} \quad \mathbf{f}_1(\xi; \pi) \geq 3 \quad \mathbf{f}_{-1}(\xi; \pi) \leq 1 \quad \mathbf{f}_2(\xi; \pi) \leq \frac{1}{2}. \end{aligned} \quad (1.8)$$

Moreover, for all $\varphi_0 \in \mathbb{T}$, $\xi \in \mathcal{O}_\varepsilon$ one has

$$\begin{cases} \mathbf{f}_1(\xi; \varphi_0) + 2\mathbf{f}_2(\xi; \varphi_0) = \xi_1 \\ \mathbf{f}_{-1}(\xi; \varphi_0) - 2\mathbf{f}_2(\xi; \varphi_0) = \xi_2 \\ \mathbf{f}_{-2}(\xi; \varphi_0) + \mathbf{f}_2(\xi; \varphi_0) = \xi_3 \end{cases} \quad (1.9)$$

while

$$\mathbf{f}_2(\cdot, \xi_1, \xi_2, \xi_3; 0) - \mathbf{f}_2(\cdot, \xi_1, \xi_2, \xi_3; \pi) \text{ is an increasing function of } \xi_0. \quad (1.10)$$

The quasi-periodic solutions $u_{\varepsilon,k}$ are analytic reducible KAM tori.

The proof of this Theorem is performed essentially in two steps, formally described in Propositions 2.1 and 2.2. We now give an informal presentation of the strategy. Note that, since the sign σ will play no role in our reasoning, we shall for simplicity set it to one in the following. We rescale $u \rightsquigarrow \varepsilon^{1/4}u$, so that the Hamiltonian \mathcal{H} assumes the form

$$\mathcal{H} = \sum_{j \in \mathbb{Z}} j^2 |u_j|^2 + \varepsilon \sum_{j_1+j_2+j_3=j_4+j_5+j_6} u_{j_1} u_{j_2} u_{j_3} \bar{u}_{j_4} \bar{u}_{j_5} \bar{u}_{j_6}, \quad (1.11)$$

while the constants of motion are unchanged. It is well known that the dynamics of (1.11) is controlled for finite but long times by the corresponding Birkhoff Hamiltonian, containing only the terms in (1.11) which Poisson commute with the quadratic part $H^{(2)} := \sum_{j \in \mathbb{Z}} j^2 |u_j|^2$, see Subsection 3. In the case of Hamiltonian 1.11 we get

$$H_{\text{BirK}} = \sum_{j \in \mathbb{Z}} j^2 |u_j|^2 + \varepsilon \sum_{\substack{j_1+j_2+j_3=j_4+j_5+j_6 \\ j_1^2+j_2^2+j_3^2=j_4^2+j_5^2+j_6^2}} u_{j_1} u_{j_2} u_{j_3} \bar{u}_{j_4} \bar{u}_{j_5} \bar{u}_{j_6}$$

which was studied in detail in [18] and [23]. A simple remark, proved for instance in [27] Propositions 1 and 2, is that while the dynamics of H_{BirK} may be very complicated still one may produce many subspaces which are invariant for the dynamics of H_{BirK} and such that the restricted dynamics is integrable.

²By the notation $O(\varepsilon^{2/3^-})$ in equation (1.7), we denote a remainder $\mathbf{r}_{\varepsilon,k}(\xi; t, x)$ belonging to the space $\ell^{a,p}$ introduced in (2.1) for some $a > 0, p > 1/2$ fixed once and for all, satisfying $\|\mathbf{r}_{\varepsilon,k}\|_{a,p} \leq C_\delta \varepsilon^{\frac{2}{3}-\delta}$ for all $\delta > 0$.

More precisely given a set $\mathcal{S} \subset \mathbb{Z}$ we define the subspace

$$U_{\mathcal{S}} := \{u \in L^2(\mathbb{T}) : u(x) = \sum_{j \in \mathcal{S}} u_j e^{ijx}\},$$

and consider the following definitions.

Definition 1.3 (Complete). *We say that a set $\mathcal{S} \subset \mathbb{Z}$ is complete if $U_{\mathcal{S}}$ is invariant for the dynamics of H_{Res} .*

Definition 1.4 (Action preserving). *A complete set $\mathcal{S} \subset \mathbb{Z}$ is said to be action preserving if all the actions $|u_j|^2$ with $j \in \mathcal{S}$ are constants of motion for the dynamics of H_{Res} restricted to $U_{\mathcal{S}}$.*

The conditions under which a given set \mathcal{S} is complete or action preserving can be rephrased more explicitly by using the structure of H_{Birk} .

Definition 1.5 (Resonance). *Given a sextuple $(j_1, \dots, j_6) \in (\mathbb{Z})^6$ we say that it is a resonance if*

$$\begin{cases} j_1 + j_2 + j_3 = j_4 + j_5 + j_6 \\ j_1^2 + j_2^2 + j_3^2 = j_4^2 + j_5^2 + j_6^2 \end{cases} \quad (1.12)$$

With this Definitions one easily sees that generic choices of \mathcal{S} lead to complete and action preserving subspaces. Indeed \mathcal{S} is complete if and only if for any quintuple $(j_1, \dots, j_5) \in \mathcal{S}^5$ there does not exist any $k \in \mathbb{Z} \setminus \mathcal{S}$ such that (j_1, \dots, j_5, k) is a resonance. Similarly \mathcal{S} is action preserving if all resonances $(j_1, \dots, j_6) \in \mathcal{S}^6$ are *trivial*, namely there exists a permutation such that $(j_1, j_2, j_3) = (j_4, j_5, j_6)$.

Note that the tori constructed in [16], [26] and [29] are all essentially supported on such generic sets \mathcal{S} . In this paper we study instead a simple choice of set \mathcal{S} which is *complete but not action preserving*. Following [18], we choose $\mathcal{S} := \{-2, -1, 1, 2\}$ and its translations, see Remark 1.6. As shown in [18], H_{Birk} restricted to \mathcal{S} is in fact still integrable (it has four degrees of freedom and three constants of motion, see Formulas 3.4 and 3.5) however it has a non-trivial structure. To study this system we restrict to a domain which is the product of a 4-cube and a 4-torus, such that all the variables in the 4-cube are constants of motion for the dynamics of $H^{(2)}$. If we take into account also the resonant terms of degree 6 we get a different qualitative behavior. Indeed after symmetry reduction it is a Hamiltonian system on a cylinder with one stable and one unstable fixed point with its separatrix, see Figure 1. Namely there is one stable and one unstable 3-torus coming from the resonant structure. We concentrate our attention on the new set of quasi periodic solutions, which are families of stable 4-tori around the new stable 3-torus. The structure of these solutions as well as their action angle variables is studied in section 4. A key point is that these tori satisfy the so-called *twist condition*, i.e. the determinant of the Hessian of H_{Birk} restricted to \mathcal{S} is not identically zero, see Lemma 4.2. In the paper [18], these tori were used to prove the existence of *beating* solutions of the (1.1) for long but finite times. This follows from the fact that the dynamics of H_{Birk} approximates the true dynamics of (1.11) for long times.

We wish to prove the persistence of *Cantor families* of such quasi periodic solutions for infinite time, as well as their linear stability by applying a KAM scheme. A main

point naturally is to control the behavior of the infinitely many *normal directions*, i.e. the Fourier coefficients u_j with $j \neq \pm 1, \pm 2$. As is well known, a crucial preliminary point is to study the reducibility of the Birkhoff system linearized at the quasi-periodic solutions. This amounts to removing the angle dependence and diagonalizing an infinite dimensional quadratic Hamiltonian.

As happens already in the case of solutions bifurcating from linear tori (see [27]), this reduction is not trivial and not of a perturbative nature but must be constructed essentially by hand. In our case (this is the content of Proposition 2.1), we proceed as follows:

1. We perform a phase shift (3.9), inspired by [27], which removes the dependence on all the angles except one. Note that such phase shifts are typical in almost-reducible finite dimensional dynamical systems, see [14]. We are left with a quadratic Hamiltonian which depends on one angle and is diagonal except for a 4×4 block.
2. We apply a phase shift (5.5) which removes the dependence on the last angle from all the terms of the Hamiltonian except for the non-diagonal 4×4 block.
3. To this last *finite dimensional* block we apply *Floquet Theorem* 5.1, in order to remove the dependence on the angle. Finally we diagonalize it by using the standard theory of quadratic Hamiltonians.

This proves item (ii) in Proposition 2.1. In order to prove (iii) we must have some control of the eigenvalues of our quadratic Hamiltonian in the normal directions, we discuss this in subsection 5.1. Finally (i), (iv) follow by a careful study of our changes of variables. Once we have proved Proposition 2.1, the NLS Hamiltonian essentially fits all the hypotheses of a KAM scheme for a PDE on the circle. In Proposition 2.2 we give a brief description of such algorithms, in order to make the paper self contained.

2 Functional setting and main Propositions

2.1 Functional setting

Given a possibly empty set $\mathcal{S} \subset \mathbb{Z}$ and setting $n = |\mathcal{S}|$ we consider the scale of complex Hilbert spaces:

$$\ell_{\mathcal{S}}^{a,p} := \left\{ \{u_k\}_{k \in \mathbb{Z} \setminus \mathcal{S}} \mid |u_0|^2 + \sum_{k \in \mathbb{Z} \setminus \mathcal{S}} |u_k|^2 e^{2a|k|} |k|^{2p} := \|u\|_{a,p}^2 \leq \infty \right\}, \quad (2.1)$$

where $a > 0$, $p > 1/2$ are fixed once and for all. When $\mathcal{S} = \emptyset$ we denote $\ell^{a,p} := \ell_{\mathcal{S}}^{a,p}$. We consider the direct product

$$E := \mathbb{C}^n \times \mathbb{C}^n \times \ell_{\mathcal{S}}^{a,p} \times \ell_{\mathcal{S}}^{a,p} \quad (2.2)$$

We endow the space E with the (s, r) -weighted norm

$$\mathbf{v} = (\varphi, y, z, \bar{z}) \in E, \quad \|\mathbf{v}\|_E := \|\mathbf{v}\|_{E,s,r} = \frac{|\varphi|_{\infty}}{s} + \frac{|y|_1}{r^2} + \frac{\|z\|_{a,p}}{r} + \frac{\|\bar{z}\|_{a,p}}{r} \quad (2.3)$$

where, $0 < s, r < 1$, and $|\varphi|_\infty := \max_{h=1,\dots,n} |\varphi_h|$, $|y|_1 := \sum_{h=1}^n |y_h|$. We shall also use the notations

$$z_j^+ = z_j, \quad z_j^- = \bar{z}_j.$$

We remark that E is a symplectic space w.r.t. the form $dy \wedge d\varphi + i \sum_k du_k \wedge d\bar{u}_k$. We also define the toroidal domain

$$D(s, r) := \mathbb{T}_s^n \times D(r) := \mathbb{T}_s^n \times B_{r^2} \times B_r \times B_r \subset E \quad (2.4)$$

where $D(r) := B_{r^2} \times B_r \times B_r$,

$$\mathbb{T}_s^n := \left\{ \varphi \in \mathbb{C}^n : \max_{h=1,\dots,n} |\operatorname{Im} \varphi_h| < s \right\}, \quad B_{r^2} := \left\{ y \in \mathbb{C}^n : |y|_1 < r^2 \right\} \quad (2.5)$$

and $B_r \subset \ell_S^{\alpha,p}$ is the open ball of radius r centered at zero. We think of \mathbb{T}^n as the n -dimensional torus $\mathbb{T}^n := 2\pi\mathbb{R}^n/\mathbb{Z}^n$.

Remark 2.1. If $n = 0$ then $D(s, r) \equiv B_r \times B_r \subset \ell^{\alpha,p} \times \ell^{\alpha,p}$.

For a vector field $X : D(s, r) \rightarrow E$, described by the formal Taylor expansion:

$$X = \sum_{\nu, i, \alpha, \beta} X_{\nu, i, \alpha, \beta}^{(\mathbf{v})} e^{i\nu \cdot \varphi} y^i z^\alpha \bar{z}^\beta \partial_{\mathbf{v}}, \quad \mathbf{v} = \varphi, y, z, \bar{z}$$

we define the *majorant* and its *norm*³:

$$\begin{aligned} MX &:= \sum_{\nu, i, \alpha, \beta} |X_{\nu, i, \alpha, \beta}^{(\mathbf{v})}| e^{s|\nu|} y^i z^\alpha \bar{z}^\beta \partial_{\mathbf{v}}, \quad \mathbf{v} = \varphi, y, z, \bar{z} \\ |X|_{s,r} &:= \sup_{(y, z, \bar{z}) \in D(s, r)} \|MX\|_{E, s, r}. \end{aligned} \quad (2.6)$$

In our algorithm we deal with Hamiltonian vector fields which depend in a Lipschitz way on some parameters $\xi \in \mathbb{R}^n$ in a compact set $\mathcal{O} \subset \mathbb{R}^n$. To handle this dependence one introduces weighted norms. Given $X : \mathcal{O} \times D(s, r) \rightarrow E$ we set

$$|X(\xi)|_{s,r,\mathcal{O}}^\gamma := \sup_{\xi \in \mathcal{O}} |X(\xi)|_{s,r} + \gamma \sup_{\eta \neq \xi \in \mathcal{O}} \frac{|X(\xi) - X(\eta)|_{s,r}}{|\xi - \eta|} \quad (2.7)$$

where γ is a parameter. Sometimes when \mathcal{O} is understood we just write $|X|_{s,r,\mathcal{O}}^\gamma = |X|_{s,r}^\gamma$.

Definition 2.1. We say that a Hamiltonian function is M -regular if its associated vector field is M -analytic i.e. has finite (2.6) norm. We define the norm of such H as

$$|H|_{s,r}^\gamma := |X_H|_{s,r}^\gamma. \quad (2.8)$$

We denote the space of real-valued functions with finite $|\cdot|_{s,r}^\gamma$ norm as $\mathcal{A}_{s,r,\mathcal{O}} \equiv \mathcal{A}_{s,r}$.

³The different weights in formula (2.3) ensure that, if $|X|_{s,r}$ is sufficiently small, then X generates a close-to-identity change of variables from $D(s/2, r/2) \rightarrow D(s, r)$.

2.2 Main Propositions

We fix once and for all⁴ $\mathcal{S} := \{-2, -1, 1, 2\}$ and $n = 4$. As explained in the introduction we first obtain a *normal form* result for our Hamiltonian (1.11).

Proposition 2.1. *There exist $\varepsilon_0, s_0, r_0 > 0$ and a compact domain $\mathcal{O}_0 \subset \mathbb{R}^4$ of positive measure such that the following holds:*

for all $\varepsilon < \varepsilon_0, s < s_0, r < r_0$ there exists an analytic function

$$\Psi : \mathcal{O}_0 \times D(s, r) \rightarrow \ell^{a,p},$$

such that for all $\xi \in \mathcal{O}_0$ the map $\Psi(\xi; \cdot) : D(s, r) \rightarrow \ell^{a,p}$ is symplectic and satisfies the following properties:

- (i) *Define $v(\xi; \varphi) = \{v_j(\xi; \varphi)\}_{j \in \mathbb{Z}} := \Psi(\xi; \varphi, y = 0, w = 0)$. Then, for all $j \in \mathcal{S}$, equations (1.8), (1.9), (1.10) hold.*
- (ii) *The NLS Hamiltonian \mathcal{H} defined in (1.11) in the new variables becomes:*

$$\mathcal{H} \circ \Psi := \mathcal{H}_{\text{NLS}} = \omega(\xi) \cdot y + \sum_{j \neq \pm 1, \pm 2} \Omega_j |w_j|^2 + \mathcal{R}, \quad \Omega_j = (j^2 + \varepsilon \Theta_j(\xi)), \quad (2.9)$$

where

$$\omega(\xi) = (0, 1, 1, 4) + \varepsilon \lambda(\xi)$$

with $\xi \rightarrow \lambda(\xi)$ an invertible analytic map and Θ_j is chosen in a finite list (in fact with 5 elements) of distinct analytic functions.

- (iii) *For all $\xi \in \mathcal{O}_0$, we have the bounds:*

$$\begin{aligned} |\lambda| + |\partial_\xi \lambda| &< M_0, \quad |\partial_\lambda \xi| < L_0, \quad |\Theta_j| + |\partial_\xi \Theta_j| < M_0, \quad \forall j \neq \pm 1, \pm 2 \\ |\lambda \cdot \ell + \Theta_i \pm \Theta_j| &> \alpha_0, \quad \forall (\ell, i, j) \neq (0, i, i) : |\ell| < 4M_0L_0 \\ |\mathcal{R}|_{s,r,\mathcal{O}_0}^1 &\leq C(\varepsilon^2 r^{-2} + \varepsilon r), \end{aligned}$$

- (iv) *\mathcal{H}_{NLS} has the constants of motion:*

$$\begin{aligned} \mathbb{L} &= y_1 + y_2 + y_3 + \sum_{j \neq \pm 1, \pm 2} |w_j|^2, \\ \mathbb{M} &= y_1 - y_2 - 2y_3 + \sum_{j=3,4} j(|w_j|^2 + |w_{-j}|^2) + \sum_{j \neq \pm 1, \pm 2, \pm 3, \pm 4} j|w_j|^2 \end{aligned}$$

Proof. The proof of this Proposition is deferred to subsection 5.2, since it requires all the structure developed in Sections 3, 4. \square

Proposition 2.2. *Consider the Hamiltonian \mathcal{H}_{NLS} in Proposition 2.1 and fix $r = \varepsilon^{1/3}$. There exist $\varepsilon_*, s_*, \gamma_*, \mathbf{C}_* > 0$ and, for all $\varepsilon \in (0, \varepsilon_*)$ and $\gamma \in (\mathbf{C}_* \varepsilon^{1/3}, \gamma_*)$, a compact domain $\mathcal{O}_\infty = \mathcal{O}_\infty(\gamma, \varepsilon) \subseteq \mathcal{O}_0$ of positive measure with $|\mathcal{O}_0 \setminus \mathcal{O}_\infty| \sim \gamma$, such that for all $s < s_0$ there exists a Lipschitz family of analytic affine symplectic maps*

$$\mathcal{O}_\infty \ni \xi \mapsto \Phi_\infty(\xi; \cdot) : D\left(\frac{s}{4}, \frac{r}{4}\right) \rightarrow D(s, r)$$

with the following properties.

⁴by Remark 1.1 we could trivially shift for any k , to $\mathcal{S} = \{k-2, k-1, k+1, k+2\}$.

(i) Φ_∞ is $O(\varepsilon^{1/3}\gamma^{-1})$ close to the identity, i.e.

$$\|\Phi_\infty(\xi; \mathbf{v}) - \mathbf{v}\|_{E,s/4,r/4} \leq \varepsilon^{1/3}\gamma^{-1} \quad \forall \mathbf{v} \in D\left(\frac{s}{4}, \frac{r}{4}\right)$$

(ii) One has

$$\mathcal{H}_{\text{NLS}} \circ \Phi_\infty := \omega^\infty \cdot y + \sum_{j \neq \pm 1, \pm 2} \Omega_j^\infty |w_j|^2 + \mathcal{R}^\infty$$

where

$$X_{\mathcal{R}^\infty} \Big|_{\substack{y=0 \\ w=0}} = 0, \quad d_{y,w} X_{\mathcal{R}^\infty} \Big|_{\substack{y=0 \\ w=0}} = 0,$$

and

$$\omega^\infty(\xi) = (0, 1, 1, 4) + \varepsilon \lambda^\infty(\xi), \quad \Omega_j^\infty(\xi) = j^2 + \varepsilon \Theta_j^\infty$$

with $\xi \rightarrow \lambda^\infty(\xi)$ an invertible Lipschitz map $|\lambda^\infty - \lambda| \sim \varepsilon^{1/3}\gamma^{-1}$, $|\Theta_j^\infty - \Theta_j| \sim \varepsilon^{1/3}\gamma^{-1}$.

Proof. The proof of this proposition is deferred to Section 6. \square

2.3 Proof of Theorem 1.1

We choose \mathcal{O}_0 as in Proposition 2.1 and $\varepsilon_0 > 0$ as the smaller between ε_0 in Proposition 2.1 and ε_* in Proposition 2.2, same for s_0 . We choose $\gamma = (-\log \varepsilon)^{-1}$ so that the hypotheses of Proposition 2.2 are fulfilled for ε small enough. Then, for all $\xi \in \mathcal{O}_\infty(\gamma, \varepsilon)$ given by Proposition 2.2, the 4-dimensional torus $\Psi \circ \Phi_\infty(\xi; \varphi, y = 0, w = 0)$ is a reducible KAM torus, by Proposition 2.2 (ii). Equations (1.7)-(1.10) hold true by Proposition 2.1 (i), by Proposition 2.2 (i), and by the definition of the norm $\|\cdot\|_{E,s,r}$ in (2.3). This gives the thesis for $k = 0$.

The fact that the thesis holds for all $k \in \mathbb{Z}$ is a trivial consequence of Remark 1.1.

3 Birkhoff Normal form

We study the Hamiltonian (1.11) which, following the definitions of Section 2, is an analytic function $: B_R \times B_R \rightarrow \mathbb{C}$ with maximal degree 6. We apply a step of Birkhoff normal form (cf. [8],[9],[1]), by which we cancel all the terms of degree 6 which do not Poisson commute with

$$H^{(2)} := \sum_{j \in \mathbb{Z}} j^2 |u_j|^2.$$

This is done by applying a well known analytic change of variables, with generating function

$$F := \varepsilon \sum_{\substack{\alpha, \beta \in (\mathbb{Z})^{\mathbb{N}}: |\alpha|=|\beta|=3 \\ \sum_k (\alpha_k - \beta_k)k=0, \sum_k (\alpha_k - \beta_k)|k|^2 \neq 0}} \binom{3}{\alpha} \binom{3}{\beta} \frac{u^\alpha \bar{u}^\beta}{\sum_k (\alpha_k - \beta_k)|k|^2}. \quad (3.1)$$

We denote the change of variables by $\Psi_{\text{Birk}} := e^{ad(F)}$ and notice that, for ε small enough, it is well defined and analytic: $B_1 \times B_1 \rightarrow B_2 \times B_2$.

By construction Ψ_{Birk} brings (1.11) to the form

$$\mathcal{H}_{\text{NLS}} = H_{\text{Birk}} + \varepsilon^2 R_{\geq 10} = \sum_{j \in \mathbb{Z}} j^2 |u_j|^2 + \varepsilon \sum_{\substack{j_1+j_2+j_3=j_4+j_5+j_6 \\ j_1^2+j_2^2+j_3^2=j_4^2+j_5^2+j_6^2}} u_{j_1} u_{j_2} u_{j_3} \bar{u}_{j_4} \bar{u}_{j_5} \bar{u}_{j_6} + \varepsilon^2 R_{\geq 10}$$

where the terms of degree 6 (which we denote by H_6) are supported on of *resonant sextuples*, see Definition 1.12. The remainder term $R_{\geq 10}$ is a M -regular analytic Hamiltonian with minimal degree 10, defined for $u \in B_1$ and with norm $|R|_1$ controlled by some ε independent constant.

Let us analyze in detail the term H_6 by studying the resonant sextuples. Among the resonant sextuples, there are the trivial ones corresponding to $\{j_1, j_2, j_3\} = \{j_4, j_5, j_6\}$. We call *action preserving* the trivial resonant sextuples (and the corresponding monomials in the Hamiltonian), while we call *effective* all the others.

We denote $\mathcal{S} = \{-2, -1, 1, 2\} \subset \mathbb{Z}$ as the set of *tangential sites*. The complementary set $\mathbb{Z} \setminus \mathcal{S}$ is the set of *normal sites*. Whenever j is a normal site, we denote $u_j = z_j$. We order the term H_6 according to the degree in the normal variables: we denote by $H_{6,k}$ the part of H_6 that is homogeneous of degree k in the z_j 's so that

$$H_6 = \sum_{k=0}^6 H_{6,k} .$$

With our choice of \mathcal{S} , the terms $H_{6,k}$ with $k \leq 2$ have a relatively simple form. In particular, the only integer solutions to (1.12) with all six elements belonging to \mathcal{S}

$$\begin{cases} (j_1, j_2, j_3) = (1, 1, -2) \\ (j_4, j_5, j_6) = (-1, -1, 2) \end{cases} \quad (3.2)$$

up to permutations of $\{j_1, j_2, j_3\}$, permutations of $\{j_4, j_5, j_6\}$ and exchange of (j_1, j_2, j_3) and (j_4, j_5, j_6) . Moreover, there are no solutions to (1.12) with five elements in \mathcal{S} and one element outside \mathcal{S} (see Lemma 2.4 of [18]) so \mathcal{S} is complete. Finally, it is easy to verify that the only integer solutions to (1.12) with four elements in \mathcal{S} and two elements outside \mathcal{S} are of the form

$$\begin{cases} (j_1, j_2, j_3) = (1, 2, -3) \\ (j_4, j_5, j_6) = (-1, -2, 3) \end{cases} \quad \begin{cases} (j_1, j_2, j_3) = (2, 2, -4) \\ (j_4, j_5, j_6) = (-2, -2, 4) \end{cases} \quad (3.3)$$

up to permutations of $\{j_1, j_2, j_3\}$, permutations of $\{j_4, j_5, j_6\}$ and exchange of (j_1, j_2, j_3) and (j_4, j_5, j_6) .

With the remarks above, we compute the terms $H_{6,k}$ with $k \leq 2$ and obtain

$$\begin{aligned} H_{6,0} &= H_{6,0}^{\text{ap}} + H_{6,0}^{\text{eff}} \\ H_{6,0}^{\text{ap}} &= 6 \left(\sum_{j \in \mathcal{S}} |u_j|^2 \right)^3 - 9 \left(\sum_{j \in \mathcal{S}} |u_j|^2 \right) \left(\sum_{j \in \mathcal{S}} |u_j|^4 \right) + 4 \left(\sum_{j \in \mathcal{S}} |u_j|^6 \right) \\ H_{6,0}^{\text{eff}} &= 9(u_1^2 u_{-2} \bar{u}_{-1} \bar{u}_2 + c.c.) \end{aligned}$$

$$\begin{aligned}
H_{6,1} &= 0 \\
H_{6,2} &= H_{6,2}^{\text{ap}} + H_{6,2}^{\text{eff}} \\
H_{6,2}^{\text{ap}} &= \left[18 \left(\sum_{j \in \mathcal{S}} |u_j|^2 \right)^2 - 9 \left(\sum_{j \in \mathcal{S}} |u_j|^4 \right) \right] \left(\sum_{j \notin \mathcal{S}} |z_j|^2 \right) \\
H_{6,2}^{\text{eff}} &= 36(u_{-1}u_{-2}\bar{u}_1\bar{u}_2z_3\bar{z}_{-3} + c.c.) + 9(u_{-2}^2\bar{u}_2^2z_4\bar{z}_{-4} + c.c.) .
\end{aligned}$$

We push the terms of degree at least three in the normal sites to the remainder, namely we write

$$\mathcal{H} = H_2 + \varepsilon H_{6,0} + \varepsilon H_{6,2} + \mathcal{R}$$

with

$$\mathcal{R} = \varepsilon H_{6,3} + \varepsilon H_{6,4} + \varepsilon H_{6,5} + \varepsilon H_{6,6} + \varepsilon^2 R_{\geq 10} .$$

Note that the constants of motion remain unchanged after the Birkhoff change of variables.

Remark 3.1. *Here and in the following we will denote by the same letter \mathcal{R} “perturbation terms” that we are for the moment ignoring, which we will bound at the end of the procedure.*

We have shown that the Birkhoff Hamiltonian restricted to the invariant subset \mathcal{S} has the form:

$$\begin{aligned}
&\sum_{j \in \mathcal{S}} j^2 |u_j|^2 + 6\varepsilon \left(\sum_{j \in \mathcal{S}} |u_j|^2 \right)^3 - 9\varepsilon \left(\sum_{j \in \mathcal{S}} |u_j|^2 \right) \left(\sum_{j \in \mathcal{S}} |u_j|^4 \right) + 4\varepsilon \sum_{j \in \mathcal{S}} |u_j|^6 \\
&\quad + 9\varepsilon (u_1^2 u_{-2} \bar{u}_{-1} \bar{u}_2 + \bar{u}_1^2 \bar{u}_{-2} u_{-1}^2 u_2)
\end{aligned} \tag{3.4}$$

with the constants of motion

$$|u_1|^2 + 2|u_2|^2, \quad |u_{-1}|^2 - 2|u_2|^2, \quad |u_{-2}|^2 + |u_2|^2 \tag{3.5}$$

For $j \in \mathcal{S}$, we first pass to the *symplectic polar coordinates* (\mathcal{I}, ϕ) namely

$$u_j = \sqrt{\mathcal{I}_j} e^{i\phi_j} \quad j \in \mathcal{S}, \tag{3.6}$$

so that the symplectic form is now $d\mathcal{I} \wedge \phi + idz \wedge d\bar{z}$. We work in the domain

$$B_{s_0, r_0} := \mathbb{I}_0 \times \mathbb{T}_{s_0}^4 \times B_{r_0} \times B_{r_0}$$

with \mathbb{I}_0 a compact domain in $(0, \infty)^4$, so that the change of variables $(\mathcal{I}, \phi, z, \bar{z}) \rightarrow u$ is well defined and analytic. One easily sees that, for ε small enough, \mathcal{R} and $\partial_{\mathcal{I}_j} \mathcal{R}$ with $j = 1, \dots, 4$, are M -regular Hamiltonians and

$$\sup_{(\mathcal{I}, \phi, z, \bar{z}) \in B_{s_0, r_0}} \left(|X_{\mathcal{R}}^{(\phi)}|_{\infty} + |X_{\mathcal{R}}^{(\mathcal{I})}|_1 + r_0^{-1} \|X_{\mathcal{R}}^{(z)}\|_{a,p} + r_0^{-1} \|X_{\mathcal{R}}^{(\bar{z})}\|_{a,p} \right) \leq C(\varepsilon r_0 + \varepsilon^2 r_0^{-1}).$$

Then we make the analytic symplectic change of coordinates $(J, \vartheta) \rightsquigarrow (\mathcal{I}, \phi)$ defined by

$$\begin{aligned}
J_0 &= \mathcal{I}_2, & J_1 &= \mathcal{I}_1 + 2\mathcal{I}_2, & J_2 &= \mathcal{I}_{-1} - 2\mathcal{I}_2, & J_3 &= \mathcal{I}_{-2} + \mathcal{I}_2 \\
\vartheta_0 &= -2\phi_1 - \phi_{-2} + 2\phi_{-1} + \phi_2 & \vartheta_1 &= \phi_1 & \vartheta_2 &= \phi_{-1} & \vartheta_3 &= \phi_{-2} .
\end{aligned} \tag{3.7}$$

Note that now J_2 is not necessarily positive. We have

$$\begin{aligned} L &= J_1 + J_2 + J_3 + \sum_{j \neq \pm 1, \pm 2} |z_j|^2 \\ M &= J_1 - J_2 - 2J_3 + \sum_{j \neq \pm 1, \pm 2} j|z_j|^2 \\ H_2 &= J_1 + J_2 + 4J_3 + \sum_{j \neq \pm 1, \pm 2} j^2|z_j|^2 \end{aligned}$$

$$\begin{aligned} H_{6,0}^{\text{ap}} &= 6(J_1 + J_2 + J_3)^3 - 9(J_1 + J_2 + J_3)(J_0^2 + (J_1 - 2J_0)^2 + (J_2 + 2J_0)^2 + (J_3 - J_0)^2) \\ &\quad + 4(J_0^3 + (J_1 - 2J_0)^3 + (J_2 + 2J_0)^3 + (J_3 - J_0)^3) \end{aligned}$$

$$H_{6,0}^{\text{eff}} = 18(J_1 - 2J_0)(J_2 + 2J_0)\sqrt{J_0(J_3 - J_0)} \cos \vartheta_0$$

$$H_{6,2}^{\text{ap}} = \left(18(J_1 + J_2 + J_3)^2 - 9(J_0^2 + (J_1 - 2J_0)^2 + (J_2 + 2J_0)^2 + (J_3 - J_0)^2) \right) \left(\sum_{j \notin \mathcal{S}} |z_j|^2 \right)$$

$$\begin{aligned} H_{6,2}^{\text{eff}} &= 72\sqrt{J_0(J_1 - 2J_0)(J_2 + 2J_0)(J_3 - J_0)} \operatorname{Re} \left(z_3 \bar{z}_{-3} e^{i(-\vartheta_0 + 3(\vartheta_2 - \vartheta_1))} \right) \\ &\quad + 18J_0(J_3 - J_0) \operatorname{Re} \left(z_4 \bar{z}_{-4} e^{i(-2\vartheta_0 + 4(\vartheta_2 - \vartheta_1))} \right). \end{aligned}$$

We now write the Hamiltonian in a more compact notation (note that $(J, \vartheta) \in \mathbb{R}^4 \times \mathbb{T}^4$):

$$\mathcal{H}_{\text{NLS}} = J_1 + J_2 + 4J_3 + \varepsilon H_{6,0}^{\text{ap}}(J) + \varepsilon H_{6,0}^{\text{eff}}(J, \vartheta_0) + \sum_{j=3,4} \tilde{H}_j + \sum_{j \neq \pm 1, \pm 2, \pm 3, \pm 4} (j^2 + \varepsilon f(J)) |z_j|^2 + \mathcal{R}$$

where

$$\tilde{H}_j = (j^2 + \varepsilon f(J))(|z_j|^2 + |z_{-j}|^2) + \varepsilon \mathcal{U}_j(J) \operatorname{Re}(e^{i(n_j \vartheta_0 + \ell_j(\vartheta_2 - \vartheta_1))} z_j \bar{z}_{-j})$$

with

$$\begin{aligned} f(J) &= 18(J_1 + J_2 + J_3)^2 - 9(J_0^2 + (J_1 - 2J_0)^2 + (J_2 + 2J_0)^2 + (J_3 - J_0)^2) \\ \mathcal{U}_3 &= 72\sqrt{J_0(J_1 - 2J_0)(J_2 + 2J_0)(J_3 - J_0)}, \\ \mathcal{U}_4 &= 18J_0(J_3 - J_0) \end{aligned} \tag{3.8}$$

and finally

$$n_3 = -1, \quad \ell_3 = 3, \quad n_4 = -2, \quad \ell_4 = 4.$$

We first make a change of variables which removes the the phases in \tilde{H}_j . We set

$$\begin{aligned} w_{-j} &= e^{-i(n_j \vartheta_0 + \ell_j(\vartheta_2 - \vartheta_1))} z_{-j}, \quad j = 3, 4 \quad w_j = z_j, \quad |j| \geq 5 \\ p &= J_0 - \sum_{j=3,4} n_j |z_{-j}|^2, \quad q = \vartheta_0, \\ K_1 &= J_1 - \sum_{j=3,4} \ell_j |z_{-j}|^2, \quad K_2 = J_2 + \sum_{j=3,4} \ell_j |z_{-j}|^2, \quad K_3 = J_3 \end{aligned} \tag{3.9}$$

and we get

$$\begin{aligned}\mathbb{L} &= K_1 + K_2 + K_3 + \sum_{j \neq \pm 1, \pm 2} |w_j|^2, \\ \mathbb{M} &= K_1 - K_2 - 2K_3 + \sum_{j=3,4} j(|w_j|^2 + |w_{-j}|^2) + \sum_{j \neq \pm 1, \pm 2, \pm 3, \pm 4} j|w_j|^2\end{aligned}$$

Now the conjugate variables are $(p, q), (K, \vartheta)$. Substituting we get the Hamiltonian

$$\mathcal{H}_{\text{NLS}} = \mathcal{H}(p, K, q) + \sum_{j=3,4} H_j(p, K, q, w_{\pm j}) + \sum_{j \neq \pm 1, \pm 2, \pm 3, \pm 4} (j^2 + \varepsilon f(p, K))|w_j|^2 + \mathcal{R}. \quad (3.10)$$

Now \mathcal{R} contains some new terms of degree at least four in $w_{\pm 3}, w_{\pm 4}$. Here

$$\mathcal{H}(p, K, q) = A(K, p) + B(K, p) \cos q$$

with

$$A(K, p) = K_1 + K_2 + 4K_3 + \varepsilon H_{6,0}^{\text{ap}}(p, K), \quad B(K, p) \cos q = \varepsilon H_{6,0}^{\text{eff}}(p, K, q) \quad (3.11)$$

while

$$H_j = (j^2 + \varepsilon f(p, K))(|w_j|^2 + |w_{-j}|^2) + \varepsilon \mathcal{U}_j(p, K) \text{Re}(w_j \bar{w}_{-j}) + \mathbf{V}_j(p, K, q) |w_{-j}|^2$$

with f, \mathcal{U}_j defined in (3.8) and

$$\mathbf{V}_j = \varepsilon \mathcal{V}_j := \ell_j \left(\partial_{K_1} \mathcal{H} - \partial_{K_2} \mathcal{H} \right) + n_j \partial_p \mathcal{H}.$$

Note that all the changes of variables are analytic provided that we require that r_0 is sufficiently small and we fix $|w|_{a,p} \leq r$ so that $|z|_{a,p} \leq r e^{6s} < r_0$. The bounds on \mathcal{R} and on its ∂_K derivatives remain the same.

4 The integrable invariant subspace

We now restrict to the invariant subspace $\{z_j = 0\}$. Our Hamiltonian is of the form

$$\mathcal{H} = A(p, K) + B(p, K) \cos q$$

with A, B defined in (3.11). The corresponding dynamical system:

$$\begin{cases} \dot{q} = \partial_p A + \partial_p B \cos q \\ \dot{p} = B \sin q \end{cases} \quad (4.1)$$

has been studied in detail in [18], here we group some facts that we will need. Let us introduce the rescaled time $\tau = \varepsilon t$, so that the system (4.1) in the rescaled time is ε -independent.

Lemma 4.1. *There exists a neighborhood $\mathcal{B} \subset \mathbb{R}^3$ of $K_\star = (4, 0, 2)$ such that the following holds:*

- (i) there exists an ε -independent analytic function $\mathfrak{p} : \mathcal{B} \rightarrow \mathbb{R}$ with $\mathfrak{p}(K_\star) = 1$ such that $q = 0, p = \mathfrak{p}(K)$ is a stable fixed point for (4.1);
- (ii) there exists an ε -independent analytic function $\mathfrak{P} : \mathcal{B} \rightarrow \mathbb{R}$ with $\mathfrak{P}(K_\star) = 1$ such that $q = \pi, p = \mathfrak{P}(K)$ is an unstable fixed point for (4.1);
- (iii) up to the 2π -periodicity in q , these are the only fixed points of the system (4.1) and the phase portrait is qualitatively the same as in Figure 1;
- (iv) for $K = K_\star$, the two homoclinic connections linking the unstable fixed points to itself intersect the axis $q = 0$ at two points with $p = p^{(1)}, p = p^{(2)}$, satisfying $p^{(1)} + p^{(2)} = 2$ and

$$p^{(1)} < \frac{1}{2}, \quad p^{(2)} > \frac{3}{2}. \quad (4.2)$$

Proof. The system is 2π -periodic in q , direct computation shows that for $K = K_\star = (4, 0, 2)$

$$\mathcal{H} = 1308 + \varepsilon \left(-270(p^2 + (2-p)^2) + 36(p^3 + (2-p)^3) + 72p^{3/2}(2-p)^{3/2} \cos q \right)$$

there is a stable fixed point at $q = 0, p = 1$ (a non degenerate maximum for the Hamiltonian) corresponding to a periodic solution and a non degenerate unstable fixed point $q = \pi, p = 1$.

Then by the Implicit Function Theorem we construct $\mathfrak{p}(K)$ by solving the equation

$$\partial_p A + \partial_p B = 0 \quad (4.3)$$

for $p = \mathfrak{p}(K)$ in a neighborhood of $K = K_\star, p = 1$. In the same way, since $B(p, K) \neq 0$ then we obtain \mathfrak{P} by solving the equation

$$\partial_p A - \partial_p B = 0. \quad (4.4)$$

Note that ε factorizes in equations (4.3), (4.4). The qualitative structure of the phase portrait follows by Morse theory. Finally, formula (4.2) is obtained by direct computation (see the evaluation of κ_\star in Section 4.1 of [18]). \square

By the previous Lemma for all $K \in \mathcal{B}$ we have an open domain delimited by the heteroclinic connections where the motion is oscillatory in p, q . Then in all of this domain minus the stable fixed point we can construct an ε -independent symplectic change of variables which conjugates the system to action angles

$$\begin{aligned} (p, K; q, \vartheta) &\xrightarrow{\Psi} (E, K; \psi, \varphi) : \quad \mathcal{H} \circ \Psi = H(E, K, \varepsilon) \\ p = p(E, K, \varphi), \quad q = q(E, K, \varphi), \quad \vartheta = \psi + \mathcal{F}(E, K, \varphi) \end{aligned} \quad (4.5)$$

note that \mathbb{L}, \mathbb{M} do not depend on p (nor on the angles) hence $\mathbb{L}(K) \circ \Psi = \mathbb{L}(K)$ and the same for \mathbb{M} .

We denote by \mathcal{D} the (open) domain of definition of H .

Lemma 4.2. *The Hamiltonian $H(E, K)$ is real-analytic in \mathcal{D} moreover there exists a proper algebraic hypersurface \mathcal{Z} such that for all $(E, K) \in \mathcal{D} \setminus \mathcal{Z}$ the determinant of the Hessian matrix $\partial_{E,K}^2 H(E, K)$ is not zero.*

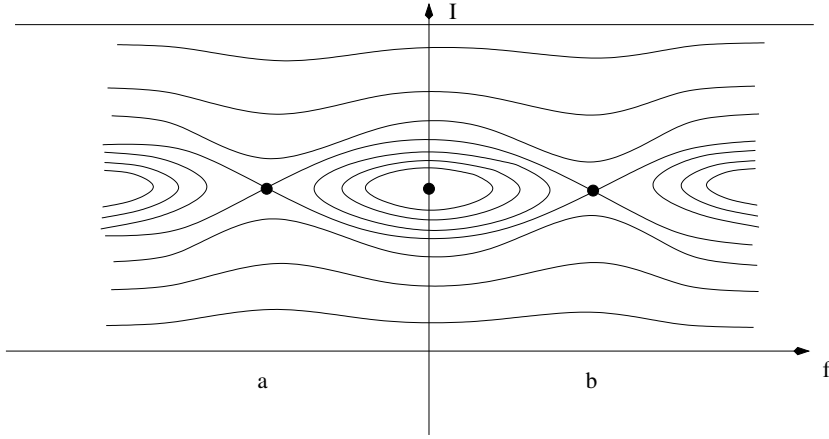


Figure 1: The phase portrait of \mathcal{H} at some value $K \in \mathcal{B}$.

Proof. We first note that the change of variables is **real** analytic in \mathcal{D} , in particular for each K in a neighborhood of K_* the p, q are analytic in E in the whole of the region delimited by the two separatrices, excluding the fixed point. We wish to compute the twist i.e. $\partial_{E,K}^2 H(E, K)$ and show that its determinant is a non-identically zero analytic function. We verify this near the stable fixed point, where the computations are explicit.

By Vey's theorem the Birkhoff normal form at the fixed point converges, namely there exists a symplectic change of variables

$$\mathbf{P} = \mathbf{P}(p - \mathbf{p}(K), q, K), \quad \mathbf{Q} = \mathbf{Q}(p - \mathbf{p}(K), q, K), \quad \phi = \vartheta + \Phi(p - \mathbf{p}(K), q, K)$$

which conjugates the Hamiltonian \mathcal{H} to $H(\mathbf{P}^2 + \mathbf{Q}^2, K)$ in a neighborhood of $\mathbf{P}, \mathbf{Q} = 0$. Note that the variables K play the rôle of parameters for \mathcal{H} . Hence we first construct symplectic changes of variables for p, q (depending parametrically on K) at the end of the procedure we complete the symplectic change of variables by adjusting the angles conjugated to K . Since the change of variable is analytic and the actions are uniquely defined we must have that $\mathbf{P}^2 + \mathbf{Q}^2 = E$ and the Birkhoff map gives an analytic extension to zero. So we may compute the E, K derivatives of H from the Birkhoff normal form at $E = 0, K = K_*$.

We start by Taylor expanding in p, q the Hamiltonian at the point $(\mathbf{p}, 0)$ up to order four. We denote $G = A + B$, see formula (3.11). Given a function $f(K, p)$ we write⁵ $f^\sharp(K) := f(K, \mathbf{p}(K))$.

$$\begin{aligned} \mathcal{H}(p, q) &= G^\sharp + \frac{1}{2}G_{pp}^\sharp(p - \mathbf{p})^2 - \frac{1}{2}B^\sharp q^2 + \frac{1}{6}G_{ppp}^\sharp(p - \mathbf{p})^3 - \frac{1}{2}B_p^\sharp(p - \mathbf{p})q^2 \\ &\quad + \frac{1}{24}B^\sharp q^4 - \frac{1}{4}B_{pp}^\sharp(p - \mathbf{p})^2 q^2 + \frac{1}{24}G_{pppp}^\sharp(p - \mathbf{p})^4 + \mathcal{O}((|q| + |p - \mathbf{p}|)^5). \end{aligned}$$

Note that $\varepsilon^{-1}(\mathcal{H} - G^\sharp)$ is ε -independent, while G^\sharp is of order $O(1)$ in ε . It follows that the Birkhoff changes of coordinates that we shall perform are all ε -independent. First we

⁵ Note that the application $f \mapsto f^\sharp$ does not commute with the derivatives. We denote $f_K^\sharp := (f_K)^\sharp$ and $\partial_K f^\sharp := \partial_K(f^\sharp)$, similarly for the derivatives with respect to p .

symmetrize the quadratic terms and translate the critical point to zero.

$$P = \lambda(p - \mathfrak{p}), \quad Q = q\lambda^{-1}, \quad \lambda^4 := -\frac{G_{pp}^\#}{B^\#}$$

(recall that $G_{pp}^\# < 0$ and $B^\# > 0$ in a neighborhood of K_\star) so, setting

$$\alpha^{(0)} = G^\#, \quad \alpha^{(2)} := \sqrt{-B^\# G_{pp}^\#},$$

we get

$$\begin{aligned} \mathcal{H} &= \alpha^{(0)} - \frac{1}{2}\alpha^{(2)}(P^2 + Q^2) + \frac{1}{6}G_{ppp}^\#\lambda^{-3}P^3 - \frac{1}{2}B_p^\#\lambda PQ^2 + \\ &+ \frac{1}{24}B^\#\lambda^4 Q^4 - \frac{1}{4}B_{pp}^\#P^2Q^2 + \frac{1}{24}G_{pppp}^\#\lambda^{-4}P^4 + \mathcal{O}(|P| + |Q|)^5 \end{aligned}$$

It is convenient to pass to complex notation in order to perform the Birkhoff normal form, let us write the Hamiltonian (neglecting higher order terms) as

$$\mathcal{H} = \alpha^{(0)} - \frac{1}{2}\alpha^{(2)}(P^2 + Q^2) + \sum_{h=3,4} \sum_{i+2j=h} \alpha_{i,j}^{(h)} P^i Q^{2j}$$

where

$$\alpha_{3,0}^{(3)} = \frac{1}{6}G_{ppp}^\#\lambda^{-3}, \quad \alpha_{1,1}^{(3)} = -\frac{1}{2}B_p^\#\lambda,$$

and

$$\alpha_{4,0}^{(4)} = \frac{1}{24}G_{pppp}^\#\lambda^{-4}, \quad \alpha_{2,1}^{(4)} = -\frac{1}{4}B_{pp}^\#, \quad \alpha_{0,2}^{(4)} = \frac{1}{24}B^\#\lambda^4.$$

we set $\sqrt{2}z = P + iQ$, so that

$$\begin{aligned} \mathcal{H} &= \alpha^{(0)} - \alpha^{(2)}|z|^2 + \sum_{h=3,4} \sum_{i+2j=h} (-1)^j \frac{\alpha_{i,j}^{(h)}}{\sqrt{2}^h} (z + \bar{z})^i (z - \bar{z})^{2j} \\ &= \alpha^{(0)} - \alpha^{(2)}|z|^2 + \sum_{h=3,4} \sum_{l+m=h} \beta_{l,m}^{(h)} z^l \bar{z}^m \end{aligned}$$

with

$$\beta_{l,m}^{(l+m)} = \frac{1}{\sqrt{2}^{l+m}} \sum_{i+2j=l+m} \sum_{\substack{0 \leq a \leq i \\ 0 \leq b \leq 2j \\ a+b=l}} (-1)^j \alpha_{i,j}^{(l+m)} \binom{i}{a} \binom{2j}{b}.$$

In particular we compute

$$\beta_{2,2}^{(4)} = \frac{3}{2}(\alpha_{0,2}^{(4)} - \alpha_{2,1}^{(4)} + \alpha_{4,0}^{(4)}) = \frac{1}{16}B^\#\lambda^4 + \frac{3}{8}B_{pp}^\# + \frac{1}{16}G_{pppp}^\#\lambda^{-4}.$$

Direct computations show that $\beta_{l,m}^{(3)}(K_\star) = 0$.

Now we need to remove the terms of order three: we perform a change of variables with generating function

$$F = \sum_{l+m=3} F_{l,m}(K) z^l \bar{z}^m$$

with $F(K_\star) = 0$ which cancels the terms of degree three. We have

$$F_{l,m} = \frac{\beta_{l,m}^{(3)}}{i\alpha^{(2)}(l-m)} \quad (4.6)$$

note that since $l + m = 3$ then $l - m$ cannot be zero.

This gives the new Hamiltonian with new terms of order four

$$\mathcal{H} = \alpha^{(0)} - \alpha^{(2)}|z|^2 + \sum_{l+m=4} \gamma_{l,m}^{(4)} z^l \bar{z}^m + h.o.t.$$

and we know that

$$\gamma_{l,m}^{(4)}(K_\star) = \beta_{l,m}^{(4)}(K_\star). \quad (4.7)$$

since $F(K_\star) = 0$. Now we remove the terms of order 4 which are not in the kernel. By definition the kernel is the part depending only on $|z|^2$ i.e. $\gamma_{2,2}^{(4)}$. This procedure may be repeated indefinitely and, by Vey's Theorem, it converges to $H(|z|^2, K)$. By construction:

$$H(|z|^2, K) = \alpha^{(0)}(K) - \alpha^{(2)}(K)|z|^2 + \gamma_{2,2}^{(4)}(K)|z|^4 + h.o.t.$$

Thus the twist matrix $\partial_{E,K}^2 H(E, K)$ evaluated at $E = 0$ is

$$M(K) := \begin{pmatrix} 2\gamma_{2,2}^{(4)} & -(\partial_K \alpha^{(2)})^T \\ -\partial_K \alpha^{(2)} & \partial_{KK} \alpha^{(0)} \end{pmatrix}.$$

Note that $\varepsilon^{-1}M(K)$ is ε -independent. We now compute the determinant of this matrix at K_\star . Note that, since we are evaluating the matrix at K_\star , we do not need to compute explicitly the corrections to the order four given by the change of variables (4.6), see formula (4.7). Direct computations show that $\det M(K_\star) \neq 0$, thus the twist condition is fulfilled outside the zero set \mathcal{Z} of an analytic function of $(E, K) \in \mathbb{C}^4$. \square

5 Normal form in the normal directions.

As we have mentioned before, $H(E, K)$ is real-analytic in \mathcal{D} and satisfies the twist conditions outside \mathcal{Z} . Now we choose a point in \mathcal{D} where we wish to prove persistence of tori by setting

$$(E, K) = \xi + y = (\xi_0, \xi_1, \xi_2, \xi_3) + (y_0, y_1, y_2, y_3), \quad (5.1)$$

where ξ are parameters and with y is in some neighborhood $B_{r,2}$ of the origin so that the action $\xi + r^2$ belongs to \mathcal{D} . We systematically use the notation

$$\mathbf{F}(\xi, \varphi) = \mathbf{F}(p(\xi_0, \xi_1 \xi_2, \xi_3, \varphi), \xi_1, \xi_2, \xi_3, q(\xi_0, \xi_1 \xi_2, \xi_3, \varphi)). \quad (5.2)$$

Now the Hamiltonian is

$$\mathcal{H}_{\text{NLS}} = \omega(\xi) \cdot y + \sum_{j=3,4} \mathcal{H}_j(\xi, \varphi, w_{\pm j}) + \sum_{j=\pm 1, \pm 2, \pm 3, \pm 4} (j^2 + \varepsilon f(\xi, \varphi)) |w_j|^2 + \mathcal{R} \quad (5.3)$$

where

$$\omega(\xi) = \partial_\xi H(\xi) =: (0, 1, 1, 4) + \varepsilon \lambda(\xi)$$

$$\mathcal{H}_j = (j^2 + \varepsilon f(\xi, \varphi))(|w_j|^2 + |w_{-j}|^2) + \varepsilon \mathcal{U}_j(\xi, \varphi) \operatorname{Re}(w_j \bar{w}_{-j}) + \varepsilon \mathcal{V}_j(\xi, \varphi) |w_{-j}|^2$$

Note that we have grouped in \mathcal{R} all the terms which are of degree higher than one in y or greater than two in w . As before the change of variables is analytic in all its entries and we have the bounds

$$|\mathcal{R}|_{s,r,\mathcal{O}}^1 \leq \mathbf{C}(\varepsilon r + \frac{\varepsilon^2}{r^2}) \quad (5.4)$$

provided that we choose for \mathcal{O} a compact domain in \mathcal{D} and choose ε, r small enough, in particular much smaller than the distance between \mathcal{O} and the boundary of \mathcal{D} .

The term $\varepsilon^2 r^{-2}$ is due to the scaling of the domain in the y -component of the Hamiltonian vector field.

This change of variables gives for the constants of motion

$$\mathbb{L} = y_1 + y_2 + y_3 + \sum_{j \neq \pm 1, \pm 2} |w_j|^2,$$

$$\mathbb{M} = y_1 - y_2 - 2y_3 + \sum_{j=3,4} j(|w_j|^2 + |w_{-j}|^2) + \sum_{j \neq \pm 1, \pm 2, \pm 3, \pm 4} j |w_j|^2$$

We wish to reduce the “normal form” above to constant coefficients. First we remove the φ dependence in w_j for $j \neq \pm 1, \pm 2, \pm 3, \pm 4$. This is done by performing the symplectic change of variables

$$y_0^{(\text{new})} = y_0 + \frac{f(\xi, \varphi) - f_0(\xi)}{\lambda_0(\xi)} \sum_{j \neq \pm 1, \pm 2} |w_j|^2, \quad w_j^{(\text{new})} = w_j \exp\left(\frac{-i \partial_\varphi^{-1} f(\xi, \varphi)}{\lambda_0(\xi)}\right), \quad (5.5)$$

here $f_0(\xi)$ denotes the average of f in φ while $\partial_\varphi^{-1} f$ is the zero average primitive of $f - f_0$. Since f is real this gives a phase shift (leaves each $|w_j|^2$ unchanged) and hence also the mass and momentum are unchanged.

Dropping the superscript $^{(\text{new})}$ the Hamiltonian is:

$$\mathcal{H}_{\text{NLS}} = \omega(\xi) \cdot y + \sum_{j=3,4} \tilde{\mathcal{H}}_j(\xi, \varphi, w_{\pm j}) + \sum_{j \neq \pm 1, \pm 2, \pm 3, \pm 4} (j^2 + \varepsilon f_0(\xi)) |w_j|^2 + \mathcal{R}$$

with

$$\tilde{\mathcal{H}}_j = (j^2 + \varepsilon f_0(\xi))(|w_j|^2 + |w_{-j}|^2) + \varepsilon \mathcal{U}_j(\xi, \varphi) \operatorname{Re}(w_j \bar{w}_{-j}) + \varepsilon \mathcal{V}_j(\xi, \varphi) |w_{-j}|^2, \quad j = 3, 4.$$

Possibly restricting to a smaller set \mathcal{O} , in order to assure that $\lambda_0(\xi)$ is bounded away from zero we have for \mathcal{R} the same bounds (5.4).

Now the dependence on φ appears only in the finite dimensional blocks given by $\tilde{\mathcal{H}}_3, \tilde{\mathcal{H}}_4$, where we reduce to constant coefficients by using Floquet’s theorem. Consider the Hamiltonian

$$H_{\text{F1}} = \varepsilon \lambda_0 y_0 + \tilde{\mathcal{H}}_3(\xi, \varphi, w_{\pm 3}) + \tilde{\mathcal{H}}_4(\xi, \varphi, w_{\pm 4})$$

Direct computation shows that the Hamiltonian flow preserves the quantities $|w_j|^2 + |w_{-j}|^2$ for $j = 3, 4$. Since the reduction procedure is the same for $j = 3, 4$, we explicitly perform it on one block.

The Hamilton equations for $\mathfrak{w}_j = (w_j, w_{-j})$ have the 2×2 block structure

$$\begin{pmatrix} \dot{\mathfrak{w}}_j \\ \dot{\bar{\mathfrak{w}}}_j \end{pmatrix} = ij^2 \begin{pmatrix} \mathfrak{w}_j \\ \bar{\mathfrak{w}}_j \end{pmatrix} + \varepsilon \begin{pmatrix} A_j(\xi, \varphi) & 0 \\ 0 & \overline{A_j(\xi, \varphi)} \end{pmatrix} \begin{pmatrix} \mathfrak{w}_j \\ \bar{\mathfrak{w}}_j \end{pmatrix}, \quad (5.6)$$

where the 2×2 matrix

$$A_j(\xi, \varphi) := i \begin{pmatrix} f_0 & \mathcal{U}_j \\ \mathcal{U}_j & f_0 + \mathcal{V}_j \end{pmatrix} \quad (5.7)$$

is skew-Hermitian and $\overline{A_j(\xi, \varphi)}$ denotes the entry-wise complex conjugate of $A_j(\xi, \varphi)$. In other words, the two Lagrangian subspaces $\mathfrak{w}_j = 0$, $\bar{\mathfrak{w}}_j = 0$ are invariant, the system is decoupled into

$$\dot{\varphi} = \varepsilon \lambda_0, \quad \dot{\mathfrak{w}}_j = (ij^2 \mathbb{1} + \varepsilon A_j(\xi, \varphi)) \mathfrak{w}_j, \quad \dot{\bar{\mathfrak{w}}}_j = (-ij^2 \mathbb{1} + \varepsilon \overline{A_j(\xi, \varphi)}) \bar{\mathfrak{w}}_j, \quad (5.8)$$

By variation of constants we may set $\mathfrak{w}_j = \mathbf{z}_j e^{ij^2 t}$ so that the equations for \mathbf{z}_j are

$$\dot{\varphi} = \varepsilon \lambda_0, \quad \dot{\mathbf{z}}_j = \varepsilon A_j(\xi, \varphi) \mathbf{z}_j, \quad \dot{\bar{\mathbf{z}}}_j = \varepsilon \overline{A_j(\xi, \varphi)} \bar{\mathbf{z}}_j, \quad (5.9)$$

which is still a Hamiltonian system with Hamiltonian

$$H_{\text{Fl}} = \varepsilon \left(\lambda_0 y_0 + \sum_{j=3,4} f_0(\xi) (|w_j|^2 + |w_{-j}|^2) + \mathcal{U}_j(\xi, \varphi) \text{Re}(w_j \bar{w}_{-j}) + \varepsilon \mathcal{V}_j(\xi, \varphi) |w_{-j}|^2 \right).$$

Let us study this system and rescale the time to $\tau = \varepsilon t$.

Theorem 5.1 (Floquet's theorem). *Consider a complex linear $n \times n$ dynamical system $\dot{x} = A(\tau)x$ where the matrix A is periodic of period T . Let $X(\tau)$ be the fundamental matrix solution. Then*

$$X(\tau) = P(\tau) e^{\tau B}$$

where the $n \times n$ matrices $P(\tau)$ and B satisfy:

- $P(\tau)$ is invertible and periodic of period T .
- B is time independent and satisfies $X(T) = e^{TB}$.

As a consequence, the dynamical system $\dot{x} = A(\tau)x$ is symplectically conjugated to the constant coefficient system $\dot{v} = Bv$ through the time-periodic change of variables $P(\tau)$.

For $j = 3, 4$, the system (5.9) satisfies all the assumptions of Theorem 5.1 and therefore we may reduce it to constant coefficients. From the proof of Floquet's theorem, it also follows that we can perform such reduction in a way that preserves the skew-Hermiticity of the matrix. In fact, the matrix B in Theorem 5.1 is constructed considering the matrix fundamental solution of

$$\begin{cases} \dot{W}_j = A_j(\tau) W_j \\ W_j(0) = \mathbb{1} \end{cases} \quad (5.10)$$

and finding a constant coefficient matrix B_j such that $e^{TB_j} = W_j(T)$. The key point is that the 2×2 skew-Hermitian matrices form the Lie algebra $\mathfrak{su}(2)$ associated to the special unitary group $SU(2)$, which is compact and connected. This implies that $W_j(t) \in SU(2)$ for all t and, in particular, $W_j(T) \in SU(2)$. Hence, using the fact that the exponential map is surjective on compact connected Lie groups, the matrix B_j can be chosen so that $B_j \in \mathfrak{su}(2)$. Thus, the time-periodic change of variables $P_j(t) := W_j(t)e^{-tB_j}$ conjugates (5.8) to the constant coefficient dynamical system (with an abuse of notation, we denote the new variables again with \mathbf{z}_j)

$$\dot{\varphi} = \varepsilon\lambda_0, \quad \dot{\mathbf{z}}_j = \varepsilon B_j(\xi)\mathbf{z}_j, \quad B_j(\xi) = i \begin{pmatrix} a_j & \bar{c}_j \\ c_j & b_j \end{pmatrix} \in \mathfrak{su}(2), \quad (5.11)$$

note that we have returned to the original time-scale.

Clearly, the change of variables $\bar{P}_j(t) := \overline{W_j(t)}e^{-t\bar{B}_j}$ conjugates the last equation in (5.9) to

$$\dot{\bar{\mathbf{z}}}_j = \varepsilon \overline{B_j(\xi)} \bar{\mathbf{z}}_j. \quad (5.12)$$

We now return to our original system (5.8). We note that the fundamental solution $W(\tau)$ of (5.8) is $W(\tau) = e^{ij^2\varepsilon^{-1}\tau}X(\tau)$ so if we apply the change of variables given by P to (5.8) we get

$$\dot{\varphi} = \varepsilon\lambda_0, \quad \dot{\mathbf{w}}_j = (ij^2 + \varepsilon B_j(\xi))\mathbf{w}_j. \quad (5.13)$$

Adding to y_0 a correction quadratic in the $w_{\pm j}$ with $j = 3, 4$, we obtain a symplectic change of variables. Thus, the Hamiltonian becomes

$$\mathcal{H}_{\text{NLS}} = \omega(\xi) \cdot y + \sum_{j=3,4} \mathcal{Q}_j(\xi, w_j, w_{-j}, \bar{w}_j, \bar{w}_{-j}) + \sum_{j \neq \pm 1, \pm 2, \pm 3, \pm 4} (j^2 + \varepsilon f_0(\xi)) |w_j|^2 + \mathcal{R} \quad (5.14)$$

with

$$\mathcal{Q}_j = j^2(|w_j|^2 + |w_{-j}|^2) + \varepsilon \{a_j |w_j|^2 + b_j |w_{-j}|^2 + \text{Re}(c_j w_j \bar{w}_{-j})\}.$$

The mass and momentum become

$$\mathbb{L} = y_1 + y_2 + y_3 + \sum_{j \neq \pm 1, \pm 2} |w_j|^2,$$

$$\mathbb{M} = y_1 - y_2 - 2y_3 + \sum_{j=3,4} j(|w_j|^2 + |w_{-j}|^2) + \sum_{j \neq \pm 1, \pm 2, \pm 3, \pm 4} j |w_j|^2$$

since the Floquet change of variables must preserve $|w_j|^2 + |w_{-j}|^2$ for $j = 3, 4$.

By the classification of quadratic Hamiltonians, since $B_j(\xi)$ is skew-hermitian there exists a symplectic change of variables, depending on $\xi \in \mathcal{O}$ smoothly, where the \mathcal{Q}_j are diagonal, i.e.

$$\mathcal{H}_{\text{NLS}} = \mathcal{N} + \mathcal{R} \quad \mathcal{N} = \omega(\xi) \cdot y + \sum_{j \neq \pm 1, \pm 2} \Omega_j |w_j|^2, \quad \Omega_j = (j^2 + \varepsilon \Theta_j(\xi)), \quad (5.15)$$

Where we have denoted by $\Theta_{\pm j}$ the eigenvalues of $-iB_j$ for $j = 3, 4$ and for $|j| > 4$ we have set $\Theta_j = f_0$. Finally \mathcal{R} satisfies bounds of the form (5.4).

5.1 Melnikov conditions

As in the paper [30], we study the action of the NLS normal form

$$\mathcal{N} := \omega(\xi) \cdot y + \sum_{j \neq \pm 1, \pm 2} \Omega_j(\xi) |w_j|^2$$

by Poisson bracket on the space of regular analytic Hamiltonians that are at most quadratic in the normal modes and commute with mass and momentum.

We only need to study the action of \mathcal{N} on monomials of the form $e^{i\ell \cdot \psi} w_h^{\sigma'} w_k^\sigma$

$$\ell = (\ell_0, \ell_1, \ell_2, \ell_3), \quad \sigma \in \{0, 1, -1\}, \quad h, k \in \mathbb{Z} \setminus \{\pm 1, \pm 2\} \quad (5.16)$$

where by definition

$$w_k^\sigma := \begin{cases} w_k & \text{if } \sigma = 1 \\ \bar{w}_k & \text{if } \sigma = -1 \\ 1 & \text{if } \sigma = 0 \end{cases} .$$

Mass and momentum conservation means that we only need to consider monomials satisfying the following constraints:

$$\eta(\ell) + \sigma' + \sigma = 0 \quad (5.17)$$

$$\pi(\ell) + \sigma' \mathbf{r}(h) + \sigma \mathbf{r}(k) = 0 \quad (5.18)$$

where

$$\eta(\ell) := \ell_1 + \ell_2 + \ell_3, \quad \pi(\ell) := \ell_1 - \ell_2 - 2\ell_3, \quad \mathbf{r}(j) := \begin{cases} |j| & \text{if } j = 3, 4 \\ j & \text{if } |j| \geq 5 \end{cases}$$

Proposition 5.1. *For each $\ell, \sigma, \sigma', h, k$ satisfying (5.16), (5.17) and (5.18) the Melnikov resonance condition*

$$\omega \cdot \ell + \sigma \Omega_h + \sigma' \Omega_k = 0 \quad (5.19)$$

defines a proper codimension one analytic manifold except in the trivial case $\ell = 0, \sigma + \sigma' = 0, h = k$.

Proof. It is sufficient to verify that none of these analytic functions are identically zero.

We first consider the second order Melnikov conditions, namely $\sigma, \sigma' \neq 0$. Since (5.19) are all affine functions of ε , these functions are identically zero if and only if

$$\pi^{(2)}(\ell) + \sigma \mathbf{r}^2(h) + \sigma' \mathbf{r}^2(k) = 0, \quad \lambda \cdot \ell + \sigma \Theta_h + \sigma' \Theta_k \equiv 0 \quad (5.20)$$

where

$$\pi^{(2)}(\ell) := \ell_1 + \ell_2 + 4\ell_3.$$

In conclusion we have the system of equations

$$\begin{cases} \ell_1 + \ell_2 + \ell_3 + \sigma + \sigma' = 0 & (5.21) \\ \ell_1 - \ell_2 - 2\ell_3 + \sigma \mathbf{r}(h) + \sigma' \mathbf{r}(k) = 0 & (5.22) \\ \ell_1 + \ell_2 + 4\ell_3 + \sigma \mathbf{r}^2(h) + \sigma' \mathbf{r}^2(k) = 0 & (5.23) \\ \lambda \cdot \ell + \sigma \Theta_h + \sigma' \Theta_k = 0 & (5.24) \end{cases}$$

Case 1: $|h|, |k| > 4$.

Recalling (5.15), we remark that if $\sigma + \sigma' = 0$, in the last equation we get $\lambda \cdot \ell = 0$. Since $\xi \rightarrow \lambda$ is generically a diffeomorphism this gives $\ell = 0$, hence $h = k$ from (5.22). If $\sigma + \sigma' = 2$ we get $\lambda \cdot \ell + 2f_0(\xi) = 0$. In order to prove that this is not an identity we only need to compute the functions $\lambda(\xi)$, $f_0(\xi)$ as the action $\xi_0 = E \rightarrow 0$ and $(\xi_1, \xi_2, \xi_3) = K \rightarrow K_\star = (4, 0, 2)$. It is easily seen that

$$\lambda \xrightarrow{\xi \rightarrow (0,4,0,2)} (\alpha^{(2)}(4, 0, 2), \partial_K \alpha^{(0)}(4, 0, 2)) = (-144\sqrt{3}, 426, 426, 498),$$

$$f_0 \xrightarrow{\xi \rightarrow (0,4,0,2)} f(1, 4, 0, 2) = 558$$

Since $-144\sqrt{3}$ is irrational while 426, 498, 558 are integer, we deduce $\ell_0 = 0$. Then we try to solve

$$\lambda(\xi) \cdot \ell + 2f_0(\xi) = 0, \quad \ell_1 + \ell_2 + \ell_3 = -2 \quad (5.25)$$

at $\xi = (0, 4, 0, 2)$ and one verifies that no integer solutions ℓ exist, since (set $\ell_1 + \ell_2 = x$, $\ell_3 = y$) the solution of the linear system

$$\begin{cases} 426x + 498y = -1116 \\ x + y = -2 \end{cases}$$

is $x = 5/3$, $y = -11/3$.

Case 2: $h, k \in \{\pm 3, \pm 4\}$.

As above, we consider the limit $\xi_0 = E \rightarrow 0$ and $(\xi_1, \xi_2, \xi_3) = K \rightarrow K_\star = (4, 0, 2)$. Recalling the notation (5.2), we compute

$$\mathcal{V}_j(\xi, \varphi) \xrightarrow{\xi \rightarrow (0,4,0,2)} \mathcal{V}_j(p = 1, K = (4, 0, 2), q = 0) = 0, \quad j = 3, 4.$$

In the same limit we get $\mathcal{U}_3 = 144$, $\mathcal{U}_4 = 18$. In this limit in (5.7) we do not have any angle dependence hence Theorem 5.1 holds trivially with $B = A$. Then we just compute the eigenvalues which are $\Theta_{\pm 3} = 558 \pm 144$, $\Theta_{\pm 4} = 558 \pm 18$. Now, since $\lambda_1, \lambda_2, \lambda_3, \Theta_h, \Theta_k$ are all integers while λ_0 is irrational (here all functions, when not specified, are evaluated at $\xi = (0, 4, 0, 2)$), equation (5.24) gives $\ell_0 = 0$. First consider the case $\sigma + \sigma' = 0$. Then equations (5.21), (5.23) imply $3\ell_3 + \sigma \mathbf{r}^2(h) + \sigma' \mathbf{r}^2(k) = 0$. Since 3 does not divide $4^2 - 3^2 = 7$, we get $|h| = |k| \in \{3, 4\}$, hence $\mathbf{r}(h) = \mathbf{r}(k)$. Equations (5.21), (5.22), (5.23) then give $\ell = 0$. The observation that $\Theta_j \neq \Theta_{-j}$ for $j = 3, 4$ implies that (5.24) with $\ell = 0$ is satisfied only if $h = k$. Now consider the case $\sigma + \sigma' = 2$. If $|h| = |k|$, then equations (5.21), (5.24) lead to the linear system

$$\begin{cases} 426x + 498y = -1116 + \rho \\ x + y = -2 \end{cases}$$

with $\rho \in \{0, \pm 36, \pm 288\}$, which has no integer solution; indeed the first equation (divided by 3) implies $x + y \equiv 0 \pmod{3}$, which is clearly incompatible with the second equation. If $|h| \neq |k|$, then equations (5.21), (5.23) imply $3\ell_3 = -23$, so there is no integer solution.

Case 3: $|h| \in \{3, 4\}$, $|k| > 4$.

Like before, we evaluate everything in the limit $\xi \rightarrow (0, 4, 0, 2)$; again, we use equation (5.24) to deduce $\ell_0 = 0$. Consider the case $\sigma + \sigma' = 2$. Then equations (5.21), (5.24) lead to the linear system (set $\ell_1 + \ell_2 = x$, $\ell_3 = y$)

$$\begin{cases} 426x + 498y = -1116 + \varrho \\ x + y = -2 \end{cases}$$

with $\varrho \in \Xi := \{\pm 18, \pm 144\}$, which has no integer solution as above. Finally, if $\sigma + \sigma' = 0$, equations (5.21), (5.24) imply $72\ell_3 = \varrho$, with $\varrho \in \Xi$. Now, if $\varrho = \pm 18$, there is no integer solution, which rules out the case $|h| = 4$. Therefore we have $|h| = 3$, $\ell_3 = \pm 2$. Using equations (5.21), (5.23), this implies $\mathbf{r}^2(k) \in \{\pm 3, \pm 15\}$, which is not possible since they are not perfect squares.

Then we consider first order Melnikov resonances of the type, without loss of generality $\sigma = -1, \sigma' = 0$.

$$\omega \cdot \ell - \Omega_h$$

cannot identically vanish. Indeed, it is enough to consider the mass conservation $\ell_1 + \ell_2 + \ell_3 = 1$ and the resonance condition $\lambda \cdot \ell = \Theta_h$. Reasoning like above, the limit $\xi \rightarrow (0, 4, 0, 2)$ leads to the linear system

$$\begin{cases} 426x + 498y = 558 + \varrho \\ x + y = 1 \end{cases}$$

with $\varrho \in \Xi = \{\pm 18, \pm 144\}$, which again has no solution, since the first equation implies $x + y \equiv 0 \pmod{3}$.

The case $\sigma = \sigma' = 0$ is completely trivial since the map $\omega \leftrightarrow \xi$ is a local diffeomorphism. This concludes the proof of the proposition. \square

Consider the Hamiltonian (5.15) in the set \mathcal{O} defined under (5.4). Let M_0, L_0 be defined by

$$|\lambda|_\infty + |\lambda|^{\text{lip}}, |f_0|_\infty + |f_0|^{\text{lip}}, \sup_{j=\pm 3, \pm 4} |\Theta_j|_\infty + |\Theta_j|^{\text{lip}} \leq M_0, \quad |(\lambda)^{-1}|^{\text{lip}} \leq L_0 \quad (5.26)$$

where given a map $f : \mathcal{O} \rightarrow \mathbb{R}^d$ we set

$$|f|_\infty := \sup_{\xi \in \mathcal{O}} \sup_{i=1, \dots, d} |f_i|, \quad |f|^{\text{lip}} := \sup_{\xi \neq \eta \in \mathcal{O}} \sup_{i=1, \dots, d} \frac{|f_i(\xi) - f_i(\eta)|}{|\xi - \eta|}.$$

Lemma 5.1. *There exists an ε -independent proper codimension one analytic manifold \mathfrak{A} so that for all $\xi \in \mathcal{O} \setminus \mathfrak{A}$ on has*

$$\omega \cdot \ell + \sigma \Omega_h + \sigma' \Omega_k \neq 0$$

for all non-trivial $\ell, \sigma, \sigma', h, k$ satisfying $|\ell| \leq 4M_0L_0$ and conditions (5.17), (5.18). Let \mathcal{O}_0 be an ε -independent compact domain in $\mathcal{O} \setminus \mathfrak{A}$. There exist constants α_0, \mathbf{R}_0 such that

$$\begin{aligned} |\lambda \cdot \ell + \sigma \Theta_h + \sigma' \Theta_k| &\geq \alpha_0 \quad \forall \sigma, \sigma' = 0, \pm 1, |\ell| < 4M_0L_0, (5.17), (5.18) \text{ hold} \\ |\mathcal{R}|_{s,r,\mathcal{O}_0}^{\gamma_0} &\leq \gamma \varepsilon \mathbf{R}_0, \quad \gamma_0 = \gamma M_0^{-1}, \end{aligned} \quad (5.27)$$

where \mathcal{R} is defined in (5.15), α_0 does not depend on ε while $\mathbf{R}_0 \sim r + \varepsilon r^{-2}$.

Proof. By Proposition 5.1 each Melnikov resonance defines a proper codimension one analytic manifold. Our first statement follows by showing that the condition $|\ell| \leq 4M_0L_0$ implies: 1. that the only Melnikov resonance surfaces which may intersect \mathcal{O} are such that

$$\pi^{(2)}(\ell) + \sigma \mathbf{r}^2(h) + \sigma' \mathbf{r}^2(k) = 0 \quad (5.28)$$

see (5.20); 2. that such resonances are only a finite number. This is due to the smallness of ε . Indeed if (5.28) does not hold, we have

$$|\lambda \cdot \ell + \sigma \Theta_i + \sigma' \Theta_j| \geq |\pi^{(2)}(\ell) + \sigma \mathbf{r}^2(h) + \sigma' \mathbf{r}^2(k)| - \varepsilon(4M_0^2L_0 + 2M_0) \geq 1 - \frac{1}{2}.$$

Then we notice that (5.28) fixes h, k inside a ball of radius $4M_0L_0$ except in the trivial case $\ell = 0, \sigma = -\sigma', h = k$. The estimates (5.27) follow trivially. \square

5.2 Proof of Proposition 2.1

We just apply all the changes of variables discussed in the previous sections. First we apply the Birkhoff change of variables generated by the Hamiltonian in (3.1), then we pass the sites $(\pm 1, \pm 2)$ to polar coordinates in (3.6) and pass to the coordinates (3.7). Next we apply our first phase shift (3.9) and then we integrate \mathcal{H} in (4.5). Now the change of variables (5.1) sets our approximately invariant torus at $y = 0, z = 0$. The NLS Hamiltonian now has the form (5.3), and the normal form depends only on one angle. In order to remove this dependence we apply the phase shift (5.5) and then Floquet's theorem 5.1. We obtain the Hamiltonian (5.14), which we diagonalize by using the standard theory of quadratic Hamiltonians. We choose \mathcal{O}_0 to be a compact domain as in Lemma 5.1. Finally, formula (1.8) is obtained as a consequence of (4.2), by possibly further restricting the set \mathcal{O}_0 . \square

6 KAM theorem

6.1 Technical set-up

We introduce a degree decomposition on $\mathcal{A}_{s,r}$ by associating to each monomial \mathbf{m} a degree $d(\mathbf{m})$ as follows

$$\mathbf{m}_{\ell,j,\alpha,\beta} = e^{i\ell \cdot \varphi} y^j z^\alpha \bar{z}^\beta \quad \rightarrow \quad d(\mathbf{m}_{\ell,j,\alpha,\beta}) = 2|j| + |\alpha| + |\beta| - 2.$$

This gives us a definition of homogeneous polynomials of degree d and a degree decomposition of analytic functions. Given $f \in \mathcal{A}_{s,r}$ we denote by $f^{(d)}$ its projection onto the homogeneous polynomials of degree d

$$f = \sum f_{\ell,j,\alpha,\beta} e^{i\ell \cdot \varphi} y^j z^\alpha \bar{z}^\beta, \quad \Pi^d f \equiv f^{(d)} = \sum_{2|j|+|\alpha|+|\beta|-2=d} f_{\ell,j,\alpha,\beta} e^{i\ell \cdot \varphi} y^j z^\alpha \bar{z}^\beta,$$

We use the same notations for $\Pi^{\leq d}$ and $\Pi^{\geq d}$. We are also interested in projections onto trigonometric polynomials in φ , we shall denote

$$\Pi_N f := \sum_{|\ell| \leq N} f_{\ell,j,\alpha,\beta} e^{i\ell \cdot \varphi} y^j z^\alpha \bar{z}^\beta, \quad \Pi_N^\perp := \mathbb{1} - \Pi_N$$

More in general given a subset of indices $\mathbb{I} \subseteq \mathbb{Z}^4 \times \mathbb{N}^4 \times \mathbb{N}^{\mathbb{Z}} \times \mathbb{N}^{\mathbb{Z}}$ we define

$$\Pi_{\mathbb{I}} f = \sum_{(\ell, j, \alpha, \beta) \in \mathbb{I}} f_{\ell, j, \alpha, \beta} e^{i\ell \cdot \varphi} y^j z^\alpha \bar{z}^\beta. \quad (6.1)$$

We denote as usual by $\{A, B\}$ the associated Poisson bracket and, if we want to stress the role of one of the two variables, we also write $\text{ad}(A)$ for the linear operator $B \mapsto \{A, B\}$.⁶ The main properties of the majorant norm (2.7), contained in [3], Lemmata 2.10, 2.15, 2.17, express the compatibility of the norm with projections and Poisson brackets, we briefly recall them in the next propositions.

Proposition 6.1. *For every $r, s > 0$ the following holds true:*

(i) *All projections $\Pi_{\mathbb{I}}$ are continuous namely*

$$|\Pi_{\mathbb{I}} h|_{s,r}^\gamma \leq |h|_{s,r}^\gamma$$

Smoothing. *One has*

$$|\Pi_{\leq N} h|_{s,r}^\gamma \leq N^{s_1} |h|_{s-s_1,r}^\gamma, \quad |\Pi_{\geq N} h|_{s,r}^\gamma \leq N^{-s_1} |h|_{s+s_1,r}^\gamma$$

(ii) *Graded Poisson algebra. Given $f, g \in \mathcal{A}_{s,r}$ for any $r' < r$ one has*

$$|\{f, g\}|_{s,r'}^\gamma \leq \left(1 - \frac{r'}{r}\right)^{-1} C(s) |f|_{s,r}^\gamma |g|_{s,r}^\gamma$$

moreover on all monomials $d(\{f, g\}) = d(f) + d(g)$.

(iii) *Partial ordering. If we have*

$$|f_{\ell, j, \alpha, \beta}| \leq |h_{\ell, j, \alpha, \beta}|, \quad f_{\ell, j, \alpha, \beta} = \overline{f_{-\ell, j, \beta, \alpha}} \quad \forall \ell, j, \alpha, \beta$$

and $h_{\ell, j, \alpha, \beta}$ is the Taylor-Fourier expansion of a function $h \in \mathcal{A}_{s,r}$ then there exists a unique function f whose Fourier expansion is $\{f_{\ell, j, \alpha, \beta}\}$ and such that

$$|f|_{s,r}^\gamma \leq |h|_{s,r}^\gamma$$

(iv) *Degree decomposition. Given a Hamiltonian $h \in \mathcal{A}_{s,r}$ which is homogeneous of degree d then $h \in \mathcal{A}_{s,R}$ for all R and one has*

$$|h|_{s,r}^\gamma \leq r^d |h|_{s,1}^\gamma$$

Proposition 6.2. *For every $r, s > 0$ the following holds true:*

(i) *Changes of variables. If $(1 - \frac{r'}{r})^{-1} |f|_{s,r}^\gamma < \varrho$ sufficiently small then its Hamiltonian vector field X_f defines a close to identity canonical change of variables \mathcal{T}_f such that:*

$$h \circ \mathcal{T}_f = e^{\{\cdot, \cdot\}} h \quad \text{satisfies} \quad |h \circ \mathcal{T}_f|_{s,r'}^\gamma \leq (1 + C\varrho) |h|_{s,r}^\gamma$$

⁶*ad* stands for *adjoint* in the language of Lie algebras.

(ii) Remainder estimates. Consider two Hamiltonians f, g with f of minimal degree \mathbf{d}_f and g of minimal degree \mathbf{d}_g , then set

$$\mathcal{P}_i(f, h) = \sum_{l=i}^{\infty} \frac{(-\text{ad}f)^l}{l!} h, \quad \text{ad}(f)h := \{h, f\} \quad (6.2)$$

then $\mathcal{P}_i(f, g)$ is of minimal degree $\mathbf{d}_f i + \mathbf{d}_g$ and we have the bounds

$$|\mathcal{P}_i(f, h)|_{s, r'}^{\gamma} \leq C(s) \left(1 - \frac{r'}{r}\right)^{-i} (|f|_{s, r}^{\gamma})^i |g|_{s, r}^{\gamma}.$$

Note that the same holds if we substitute in (6.2) the sequence $\{\frac{1}{l!}\}$ with any sequence $\{b_l\}$ such that $\forall l$ one has $|b_l| \leq \frac{1}{l!}$.

Note that \mathcal{H} defined in (1.11) is M -regular⁷ for all r .

6.2 The iterative algorithm

Definition 6.1. We first define the diagonal Hamiltonians which are our normal forms.

$$\mathcal{F}_{\text{ker}} := \{h = q \cdot y + \sum_{j \neq \pm 1, \pm 2} Q_j |z_j|^2, \quad \text{with } q \in \mathbb{R}^4, Q_j \in \mathbb{R}\},$$

Given a M -analytic Hamiltonian H we denote by $\Pi_{\text{ker}} H$ the projection of H onto \mathcal{F}_{ker} (this is a projection of the form (6.1)). Now we define $\mathcal{F}_{s, r}$ be the subspace of M -analytic Hamiltonians H such that $H - \Pi_{\text{ker}} H \in \mathcal{A}_{s, r}$ is a M -regular Hamiltonian and Poisson commutes with \mathbb{L}, \mathbb{M} . Finally let

$$\mathcal{F}_{\text{rg}} = (\mathbb{1} - \Pi_{\text{ker}}) \Pi^{\leq 0} \mathcal{F}_{s, r}.$$

so that

$$\mathcal{F}_{s, r} = \mathcal{F}_{\text{ker}} \oplus \mathcal{F}_{\text{rg}} \oplus \mathcal{F}_{s, r}^{>0}. \quad (6.3)$$

Note that $\mathcal{F}_{s, r}$ is naturally decomposed in terms of subspaces of growing degree.

By construction the NLS Hamiltonian

$$\mathcal{H}_{\text{NLS}} = \omega(\xi) \cdot y + \sum_{j \neq \pm 1, \pm 2} \Omega_j(\xi) |w_j|^2 + \mathcal{R} = \mathcal{N} + \mathcal{R} \in \mathcal{F}_{s, r}, \quad \forall r < r_0, s < s_0 \quad (6.4)$$

Remark 6.1 (The goal). By definition, the normal form \mathcal{N} is in \mathcal{F}_{ker} . In general, the condition for a Hamiltonian $H = N + P$, $N = \Pi_{\text{ker}} H$ to have KAM tori is $\Pi_{\text{rg}} H = \Pi_{\text{rg}} P = 0$. Thus, our goal is to find a symplectic transformation Φ_{∞} such that

$$\Pi_{\text{rg}}(H \circ \Phi_{\infty}) = 0.$$

⁷it is well known that the NLS is locally well posed under much weaker regularity conditions. This is not the purpose of the present paper.

The strategy is to construct this as a limit of a *quadratic Nash-Moser algorithm*. We will denote:

$$N := \Pi_{\ker}(H), \quad P_{\text{rg}} := \Pi_{\text{rg}}(H), \quad P^{>0} := \Pi^{>0}(H), \quad (6.5)$$

so that $H = N + P_{\text{rg}} + P^{>0}$.

By the Poisson algebra property (ii) of Proposition 6.1, if A has degree $d > 0$ then $\text{ad}(A)$ is strictly lower triangular on $\mathcal{F}_{s,r}$ w.r.t. the degree decomposition. On the other hand, if A has degree $d = 0$, then $\text{ad}(A)$ is block diagonal w.r.t. the degree decomposition.

We start with the NLS Hamiltonian (5.15) which we denote by $H_0 := N_0 + P_{\text{rg},0} + P_0^{>0}$, where $N_0 = \mathcal{N}$, $P_{\text{rg},0} + P_0^{>0} = \mathcal{R}$, so that

$$P_{\text{rg},0} \sim \varepsilon^2, \quad P_0^{>0} \sim \varepsilon r$$

are appropriately small.

We wish to find a convergent sequence of changes of variables

$$\Phi_{m+1} := e^{\text{ad}(F_m)} \circ \Phi_m, \quad (6.6)$$

dependent on a sequence K_m of ultraviolet cuts, so that at each step $H_{m+1} = \Phi_{m+1}(H_0) = N_{m+1} + P_{\text{rg},m+1}^{<0} + P_{m+1}^{>0}$ is such that N_m stays close to \mathcal{N} , $P_m^{>0}$ stays bounded while $P_{\text{rg},m}$ converges to zero (super-exponentially).

At a purely formal level we would like $P_{\text{rg},m+1}$ to be *quadratic* w.r.t. $P_{\text{rg},m}$. The generating function $F_m \in \mathcal{F}_{\text{rg}, \leq K_{m+1}} := \Pi_{\leq K_{m+1}} \mathcal{F}_{\text{rg}}$ is fixed by solving the homological equation

$$\{N_m, F_m\} + \Pi_{\text{rg},m} \{P_m^{>0}, F_m\} = \Pi_{\leq K_{m+1}} P_{\text{rg},m}, \quad \Pi_{\text{rg},m} := \Pi_{\text{rg}, \leq K_{m+1}} \quad (6.7)$$

which uniquely determines F_m as a *linear* function of $P_{\text{rg},m}$ provided that the linear operator:

$$\mathfrak{L}_m := \text{ad}(N_m) + \Pi_{\text{rg},m} \text{ad}(P_m^{>0}) = \text{ad}(N_m) + \Pi_{\text{rg},m} \text{ad}(P_m^1) + \Pi_{\text{rg},m} \text{ad}(P_m^2),$$

is invertible on $\mathcal{F}_{\text{rg}, \leq K_{m+1}}$ (clearly we also need some quantitative control on the inverse).

Remark 6.2. *On $\mathcal{F}_{s,r}$ the operators $\text{ad}(N_m)$, $\Pi_{\text{rg},m} \text{ad}(P_m^1)$, $\Pi_{\text{rg},m} \text{ad}(P_m^2)$ have respectively degree 0,1,2 so it should be clear that \mathfrak{L}_m is invertible if and only if $\text{ad}(N_m)$ is invertible and in this case one inverts*

$$\mathfrak{L}_m = \text{ad}(N_m) \left(1 + \text{ad}(N_m)^{-1} \Pi_{\text{rg},m} \text{ad}(P_m^{>0}) \right)$$

by inverting the second factor. This is of the form $1+A$ with A a sum of two linear operators of degree 1,2 respectively, so $A^3 = 0$ and we invert $1+A$ with the 3 term Neumann series $1 - A + A^2$.

We now justify our choice of F_m by computing one *KAM step*, for notational convenience we drop the pedex m in H_m and substitute H_{m+1} with H_+ (same notation for the other quantities involved).

Let us compute $H_+ := e^{\text{ad}(F)}H$. First split the operator $e^{\text{ad}(F)} = 1 + \text{ad}(F) + E_F$, by definition E_F is quadratic in F and hence quadratic in P_{rg} . Regarding the term

$$(1 + \text{ad}(F))(N + P_{\text{rg}} + P^{>0}) = N + P_{\text{rg}} + P^{>0} - \{N + P^{>0}, F\} + \{F, P_{\text{rg}}\}$$

we first notice that, since F is linear w.r.t. P_{rg} , then the last summand is quadratic in P_{rg} . Moreover, since F solves the homological equation, we have

$$P_{\text{rg}} - \{N + P^{>0}, F\} = (\Pi_{\text{ker}} + \Pi_{>0} + \Pi_{>K}\Pi_{\text{rg}})\{P^{>0}, F\} + \Pi_{>K}P_{\text{rg}}.$$

Then we deduce that

$$\begin{aligned} \Pi_{\text{ker}}e^{\text{ad}(F)}H &:= N_+ = N + \Pi_{\text{ker}}\{P^{>0}, F\} + \Pi_{\text{ker}}Q(P_{\text{rg}}), \\ \Pi_{\text{rg}}e^{\text{ad}(F)}H &:= P_+^{\leq 0} = \Pi_{>K}(P_{\text{rg}} + \Pi_{\text{rg}}\{P^{>0}, F\}) + \Pi_{\text{rg}}Q(P_{\text{rg}}), \\ \Pi_{>0}e^{\text{ad}(F)}H &:= P_+^{>0} = P^{>0} + \Pi_{>0}\{P^{>0}, F\} + \Pi_{>0}Q(P_{\text{rg}}), \end{aligned}$$

where $Q(P_{\text{rg}})$ is quadratic in P_{rg} and collects the terms from $E_F(H)$ and $\{F, P_{\text{rg}}\}$. We remark that $\{F, N\}$, and therefore $E_F(N)$, are fixed by the homological equation (6.7).

We now introduce some parameters that control H_m . Recall that by definition of \mathcal{F}_{ker}

$$N_m = \omega^{(m)} \cdot y + \sum_{j \neq \pm 1, \pm 2} \Omega_j^{(m)} |w_j|^2,$$

and set

$$\omega^{(m)} = (0, 1, 1, 4) + \varepsilon \lambda^{(m)}, \quad \Omega_j^{(m)} = j^2 + \varepsilon \Theta_j^{(m)}.$$

Fix a small $\gamma > 0$. Let \mathcal{O}_m be a positive measure Cantor set, assume that $\xi \rightarrow \lambda^{(m)}(\xi)$ is invertible and let $L_m, M_m, \mathbf{R}_m, \alpha_m \geq 0$ be such that

$$\begin{aligned} |\lambda^{(m)}| + |\lambda^{(m)}|^{\text{lip}} &\leq M_m, \quad |\Theta_j^{(m)}| + |\Theta_j^{(m)}|^{\text{lip}} \leq M_m, \quad |(\lambda^{(m)})^{-1}|^{\text{lip}} \leq L_m \\ |\lambda^{(m)} \cdot \ell + \Theta_i^{(m)} + \sigma \Theta_j^{(m)}| &\geq \alpha_m \quad \forall \sigma = 0, \pm 1, \forall (\ell, i, j) \neq (0, i, i) : |\ell| < 4M_0L_0 \\ |H_m - N_m|_{s,r,\mathcal{O}_m}^{\gamma_m} &\leq \gamma \varepsilon \mathbf{R}_m, \quad \gamma_m = \gamma M_m^{-1}. \end{aligned} \quad (6.8)$$

Note that our conditions are fulfilled by the Hamiltonian (5.15) for $m = 0$.

Definition 6.2. *We say that a positive parameter $b = \{b_m\}_{m \in \mathbb{N}}$ is telescopic if for each step m we have $b_0/2 < b_m < \frac{3}{2}b_0$ (usually b_m is either an increasing or a decreasing sequence).*

Note that a sufficient condition is that

$$\sum_{m=0}^{\infty} |b_{m+1} - b_m| \leq \frac{b_0}{2}. \quad (6.9)$$

We choose as in [29]:

$$r_{m+1} = (1 - 2^{-m-3})r_m, \quad s_{m+1} = (1 - 2^{-m-3})s_m, \quad K_m = 4^m K_0. \quad (6.10)$$

We choose the set \mathcal{O}_{m+1} as

$$\mathcal{O}_{m+1} := \{\xi \in \mathcal{O}_m : |\omega^{(m)} \cdot \ell + \sigma \Omega_i^{(m)} + \sigma' \Omega_j^{(m)}| \geq \varepsilon \gamma K_m^{-\tau}, \sigma, \sigma' = 0, \pm 1, |\ell| \leq K_m, \quad (6.11)$$

$$\eta(\ell) + \sigma + \sigma' = 0, \quad \pi(\ell) + \sigma \mathbf{r}(i) + \sigma' \mathbf{r}(j) = 0, \quad (\sigma \sigma', \ell, i) \neq (-1, 0, j)\}.$$

Lemma 6.1. *For all $\xi \in \mathcal{O}_{m+1}$, (6.7) admits a unique solution and moreover*

$$\|F_m\|_{s'_m, r'_m, \mathcal{O}_{m+1}}^{\gamma_m} \leq (\varepsilon\gamma)^{-1} (\mathbf{R}_m K_m^{2\tau})^3 \|P_{\text{rg}, m}\|_{s_m, r_m, \mathcal{O}_m}^{\gamma_m}, \quad (6.12)$$

where $s'_m = (s_m + s_{m+1})/2$, $r'_m = (r_m + r_{m+1})/2$, and $\gamma_m = \gamma M_m^{-1}$.

Proof. This is a subcase of Lemma 14 of [29] or Proposition 2.31 of [12]. We fix $\mathfrak{D} := \text{ad}(N)$. We first notice that (6.11) and item (iii) of Proposition 6.1 imply that for any $g \in \mathcal{F}_{\text{rg}, m}$ one has

$$|\mathfrak{D}^{-1}g|_{s, r, \mathcal{O}_{m+1}}^{\gamma_m} \leq 3(\varepsilon\gamma)^{-1} K_m^{2\tau} |g|_{s, r, \mathcal{O}_{m+1}}^{\gamma_m}.$$

Following Remark 6.2 we set $A = \mathfrak{D}^{-1}\Pi_{\text{rg}, m}\text{ad}(P_m^{>0})$. Using (6.8) and Propositions 6.1, 6.2 we get

$$|A^j g|_{s', r', \mathcal{O}_{m+1}}^{\gamma_m} \leq (K_m^{2\tau} \mathbf{R}_m)^j |g|_{s, r, \mathcal{O}_{m+1}}^{\gamma_m}, \quad j = 1, 2.$$

The thesis follows by Remark 6.2. \square

Proposition 6.3. *For K_0, τ large and for all ε, r such that $r + \varepsilon r^{-2} \leq \gamma K_0^{-6\tau-1}$, at each step m , in the set \mathcal{O}_m , one has the estimate*

$$\|P_{\text{rg}, m}\|_{s_m, r_m, \mathcal{O}_m}^{\gamma_m} \leq (r + \varepsilon r^{-2}) e^{-\frac{3}{2}m}. \quad (6.13)$$

Moreover the constants \mathbf{R}, L, M, α are telescopic. Finally our algebraic algorithm converges on the set $\cap_m \mathcal{O}_m$, and we obtain a change of variables $\Phi_\infty := \lim_{m \rightarrow \infty} \Phi_m$. The corresponding Hamiltonian H_∞ , has reducible KAM tori.

Proof. These estimates are all standard. The condition on ε, r implies that $P_0 \sim \varepsilon r + \varepsilon^2 r^{-2}$ is small w.r.t. the size of the divisor $\sim \varepsilon\gamma$. By induction, suppose we have reached some step m and proved the corresponding estimates. Using (6.12) and items (i-ii) of Proposition 6.2, one proves that F_m defines a symplectic change of variables with $H_{m+1} := e^{\text{ad}(F_m)} H_m$ well defined in the domain $D(s_{m+1}, r_{m+1})$. The corrections of the parameters \mathbf{R}, L, M, α are obtained by item (i) Proposition 6.1 and item (ii) of Proposition 6.2 using the fact that the coordinate change is very close to the identity due to the super-exponential decay of the norm of F . This implies easily the telescopic nature of the parameters used, as $|b_{m+1} - b_m| < \text{const } b_0 e^{-(3/2)^m}$. Then the algorithm converges. \square

6.3 Proof of Proposition 2.2.

The change of variables of item (i) is $\Phi_\infty := \lim_{m \rightarrow \infty} \Phi_m$ defined in (6.6) and with F_m defined in (6.7). The estimates on $\Phi_\infty - \mathbb{1}$ follow from Lemma 6.1, Proposition 6.3 and Proposition 6.2, provided that we choose $r = \varepsilon^{1/3}$.

Item (ii) follows directly from Proposition 6.3 and Definition 1.2.

In order to complete the proof of the KAM theorem we only need to show that the set $\mathcal{O}_\infty := \cap_{m=0}^\infty \mathcal{O}_m$ has positive measure. Let us denote

$$\mathfrak{R}_{\sigma, \sigma', \ell, i, j}^{(m)} := \{\xi \in \mathcal{O}_m : |\omega^{(m)} \cdot \ell + \sigma \Omega_i^{(m)} + \sigma' \Omega_j^{(m)}| < \varepsilon \gamma K_m^{-\tau}\}$$

so that

$$\mathcal{O}_{m+1} = \mathcal{O}_m \setminus \cup_{\sigma, \sigma', \ell, i, j}^* \mathfrak{R}_{\sigma, \sigma', \ell, i, j}^{(m)}$$

where the \cup^* is the union of all the sets such that

$$|\ell| < K_m, \quad \eta(\ell) + \sigma + \sigma' = 0, \quad \pi(\ell) + \sigma \mathbf{r}(i) + \sigma' \mathbf{r}(j) = 0, \quad (\sigma\sigma', \ell, i) \neq (-1, 0, j).$$

We first consider the case $\sigma\sigma' \neq 0$, corresponding to *second Melnikov conditions*. We claim that for any fixed $\ell \in \mathbb{Z}^4$, the union $\cup_{\sigma, \sigma', \ell, i, j}^*$ is finite, with an upper bound on the cardinality depending on $|\ell|$.

By estimates (6.8) and the fact that α_m is telescopic we deduce that, if we choose K_0 so that $\alpha_0 > 2K_0^{-\tau}$, we have $\mathfrak{R}_{\sigma, \sigma', \ell, i, j}^{(m)} = \emptyset$ for all $|\ell| \leq 6L_0M_0$.

Denote $\mathbf{v} := (0, 1, 1, 4)$. If $|\sigma i^2 + \sigma' j^2| > 2|\mathbf{v}||\ell|$ we have $\mathfrak{R}_{\sigma, \sigma', \ell, i, j}^{(m)} = \emptyset$. Since

$$|\sigma i^2 + \sigma' j^2| > ||i| - |j||(|i| + |j|),$$

this condition ensures that \cup^* runs only over

$$|i|, |j| \leq \text{const}|\ell|$$

unless one has $\sigma = -\sigma'$ and $|i| = |j|$.

In this last case we use momentum conservation, if $i = j$ we get $\ell = 0$ and this is explicitly excluded in our \cup^* . If $i = -j$ with $i = \pm 3, \pm 4$ again we get $\ell = 0$ so by our previous argument $\mathfrak{R}_{\sigma, \sigma', \ell, i, j}^{(m)} = \emptyset$. Finally if $i = -j \neq \pm 3, \pm 4$ we obtain

$$2|i| \leq |\pi(\ell)|,$$

which obviously implies $|i| \leq \text{const}|\ell|$.

Let us now give a bound for the measure of a single set $\mathfrak{R}_{\sigma, \sigma', \ell, i, j}^{(m)}$. By construction all the maps $\xi \rightarrow \lambda^{(m)}$ are invertible with Lipschitz inverse, so we change the variables to $\lambda \in \lambda^{(m)}(\mathcal{O}_m)$, note that this set is contained in $[-2M_0, 2M_0]^4$. We have

$$|\mathfrak{R}_{\sigma, \sigma', \ell, i, j}^{(m)}| \leq L_m^4 \left| \left\{ \lambda \in \lambda^{(m)}(\mathcal{O}_m) : |(\mathbf{v} + \varepsilon\lambda) \cdot \ell + \sigma \tilde{\Omega}_i^{(m)}(\lambda) + \sigma' \tilde{\Omega}_j^{(m)}(\lambda)| < \varepsilon\gamma K_m^{-\tau} \right\} \right|$$

We note that $|\tilde{\Omega}_i^{(m)}(\lambda)|^{\text{lip}} \leq \varepsilon L_m M_m \leq \varepsilon|\ell|/3$, since by hypothesis $|\ell| > 6M_0L_0$. Now we can introduce an orthogonal basis for \mathbb{R}^4 where the first basis vector is parallel to ℓ , in this basis

$$|\partial_{\lambda_1}(\mathbf{v} + \varepsilon\lambda) \cdot \ell| = \varepsilon|\ell|,$$

so we may estimate the measure of the resonant set as

$$|\mathfrak{R}_{\sigma, \sigma', \ell, i, j}^{(m)}| \leq \text{const} \gamma L_0 (L_0 M_0)^3 |\ell|^{-1} K_m^{-\tau}$$

Hence, we estimate

$$\begin{aligned} \left| \cup_m \cup_{\sigma\sigma' \neq 0}^* \mathfrak{R}_{\sigma, \sigma', \ell, i, j}^{(m)} \right| &\leq \left| \bigcup_m \bigcup_{\substack{|i| \leq c|\ell|, |\ell| \leq K_m, \sigma\sigma' \neq 0 \\ \sigma \mathbf{r}(i) + \sigma' \mathbf{r}(j) = \pi(\ell)}} \mathfrak{R}_{\sigma, \sigma', \ell, i, j}^{(m)} \right| \\ &\leq \text{const} \gamma L_0 (L_0 M_0)^3 \sum_{m=0}^{\infty} K_m^{-\tau+4} \leq \text{const} \gamma L_0 (L_0 M_0)^3 \end{aligned}$$

provided that $\tau > 4$.

In the remaining cases with $\sigma\sigma' = 0$ (Diophantine condition on $\lambda^{(m)}$ and *first Melnikov conditions*), one similarly obtains the corresponding estimate

$$\left| \bigcup_m \bigcup_{\sigma\sigma'=0} \mathfrak{R}_{\sigma,\sigma',\ell,i,j}^{(m)} \right| \leq \text{const } \gamma L_0 (L_0 M_0)^3$$

(actually, when $\sigma\sigma' = 0$, this estimate is simpler to obtain and completely standard). Hence

$$|\mathcal{O}_0 \setminus \mathcal{O}_\infty| \leq \left| \bigcup_m \bigcup_{\sigma,\sigma',\ell,i,j} \mathfrak{R}_{\sigma,\sigma',\ell,i,j}^{(m)} \right| \leq \text{const } \gamma L_0 (L_0 M_0)^3.$$

□

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