# Pitfalls in GPR Data Interpretation: False Reflectors Detected in Lunar Radar Cross Sections by Chang'e-3

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Abstract—Chang'e-3 (CE-3) has been the first spacecraft to soft land on the moon since the Soviet Union's Luna 24 in 1976. The spacecraft arrived at Mare Imbrium on December 14, 2013, and the same day, Yutu lunar rover separated from lander to start its exploration of the surface and the subsurface around the landing site. The rover was equipped, among other instruments, with two lunar penetrating radar systems having a working frequency of 60 and 500 MHz. The radars acquired data for about two weeks while the rover was slowly moving along a path of about 114 m. At navigation point N0209, the rover got stacked into the lunar soil and after that only data at a fixed position could be collected. The low-frequency radar data have been analyzed by different authors and published in two different papers, which reported totally controversial interpretations of the radar cross sections. This paper is devoted to resolve such controversy by carefully analyzing and comparing the data collected on the moon by Yutu rover and on earth by a prototype of LPR mounted onboard a model of the CE-3 lunar rover. Such analysis demonstrates that the deep radar features previously ascribed to the lunar shallow stratigraphy are not real reflectors, rather they are signal artifacts probably generated by the system and its electromagnetic interaction with the metallic rover.

Index Terms—Chang'e-3 (CE-3), ground penetrating radar (GPR), moon, noise, signal analysis.

#### I. INTRODUCTION

PLANETARY subsurfaces are particularly suitable environments for ground penetrating radar (GPR) investigations as they are usually dry and cold, allowing good penetration and low attenuation of the radar signals. GPR is a well-established and mature technology for earth applications [1], but is still in its infancy in planetary exploration. Indeed, since the early days of GPR development [2], only

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TABLE I
MAIN PLANETARY GPR EQUIPMENTS

Instrument	Mission (Year)	Object	Features
ALSE	Apollo 17	Moon	Orbiter;
	(1972)		5-5.5MHz
LRS	KAGUYA	Moon	Orbiter;
	(2007)		4-6MHz
			1MHz,15MHz
			Optional
LPR	Chang'e-3	Moon	Rover;
	(2013)		40MHz-80MHz;
			250MHz-750MHz
CONSERT	Rosetta	67P/C-G	Orbiter-Lander;
	(2004)	comet	90 MHz
MARSIS	Mars Express	Mars	Orbiter;
	(2003)		1.3-2.3,2.5-3.5,
	•		3.5-4.5,4.5-5.5MHz
SHARAD	MRO	Mars	Orbiter;
	(2005)		15MHz-25MHz

few space missions have been equipped with a subsurface radar instrument (see Table I), even though the interest for this type of geophysical technique has grown quite rapidly in time. The role of orbiting subsurface radars has been fundamental in the past [3] and still ongoing missions to the moon and mars [4]–[6] and will be essential to proof the existence of liquid water inside the icy crusts of the Jovian moons [7]. So far, among the successful missions reported in Table I, Chang'e-3 (CE-3) lunar penetrating radar (LPR) represents the first and only radar instrument employed onboard a rover [8], even though in the near future various rovers equipped with a GPR are expected to land on mars [9]–[11] and on the moon [12], [13].

The concept of radars mounted on a moving vehicle with the antennas operating near the surface is particularly appealing as they can provide high-resolution electromagnetic imaging of the subsurface stratigraphy at different depth, depending on the antenna frequency employed. Such imaging could be used to choose the best location for drilling [11] or could be processed to extract quantitative information on the electromagnetic properties of the subsurface to better constrain orbiting radar data (ground truth) or geological interpretations [14].

One of the major drawbacks in using near-surface antennas is the clutter generated by objects on or above the ground, a problem well known in common applications of GPR on earth and pointed out by several authors as early as the 1990s [15]. In some circumstances, especially when the direction of the antenna dipoles is perpendicular to the objects [16], the radar cross sections can be affected by strong events coming from objects located above the survey line such as metal objects, trees, boulders, and walls; these events are usually larger than those coming from subsurface reflectors because radar signal in the ground attenuates exponentially whereas signals propagating in the air attenuate only geometrically [15]. Furthermore, as the dielectric contrast between air and surface objects is strong, even far distance targets can produce overlapping events that can be interpreted as real reflectors [17]. Particular attention should be paid at subparallel events coming from above ground targets because they are difficult to be recognized (see [16], [17]) and can lead to erroneous geological interpretation.

In planetary exploration, the types of objects present above ground are limited to rocks, boulders, and topographic reliefs; however, the main problem is represented by the interaction between the emitted GPR signals and the metallic rover [18], [19]. The interference from the body of the rover can be quite strong if the antennas are unshielded and/or elevated above the ground and can introduce artifacts that can mask the subsurface events or "create" false reflectors in the radar cross section [19]. The main source of such artifacts is the energy reflected directly from the rover or after reflection from the surface, which can originate ring-down periodic signals that can be interpreted as stratigraphy or multiple reflections [20]. Moreover, further sources of clutter could derive from spurious signals generated by electronic systems and connections, wheel motion, or other instruments onboard the rover. Note that such artifacts could be filtered out by applying different processing techniques (see [21]).

The above mentioned artifacts and clutter are well visible on LPR data, especially on those collected with the 60-MHz dipolar antennas which were mounted above ground on the back of Yutu rover [22]. In fact, the top part of the radar sections collected at this frequency on the moon is systematically affected by a large ringing due to the antennasrover coupling, whereas in the lower part, the radar signal is quite weak and difficult to be interpreted. For this reason, so far, most of the work on LPR radar data has been focused on 500 MHz [23]-[26] and only two articles attempted to interpret some deep reflectors detectable on 60-MHz data as stratigraphic interfaces, obtaining totally controversial results [27], [28]. This paper is devoted to resolve such controversy by a careful analysis and comparison of data collected on the moon by Yutu rover and on earth by a prototype of LPR mounted onboard a model of the CE-3 lunar rover [29]. Our analysis shows that the deep reflectors cannot be ascribed to real geological interfaces as clamed in [27] and [28], rather they are artifacts probably introduced by the overall system. This paper is organized as follows. In Section II, the LPR and moon operations are presented. In Section III, LPR calibration on earth and previous studies is listed. Based on the radar equation, the reflector detectability of LPR is calculated, which is aimed to verify whether the radar can reach those depths, which is shown

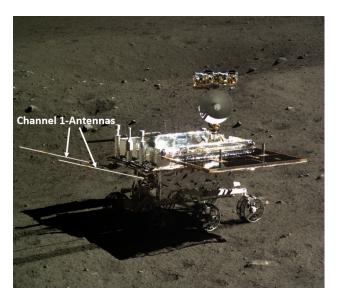


Fig. 1. Photograph of Yutu rover with channel 1 antennas (white arrows). The photograph was taken by the terrain camera onboard the lander when the rover was at navigation point N0104.

in Section IV. Two detailed data analysis, effect of noise on reflectors detectability, and time–frequency analysis applying S-transform are conducted in Section V. Section VI is the discussion. Finally, conclusions are drawn in Section VII.

#### II. LPR AND MOON OPERATIONS

CE-3 mission is part of the second phase of China Lunar Exploration Project, which is aimed at exploring the moon with rovers and landers, and follows the first phase performed by Chang'e-1 and Chang'e-2 orbiting spacecrafts [30]. CE-3 mission was successfully launched from the Xichang Satellite Launching Center on December 2, 2013, and landed on the northern Mare Imbrium after 12 days. It was the first spacecraft to soft land on the moon since the Soviet Union's Luna 24 in 1976. The same day of landing (December 14, 2013 at 20.35 UTC), Yutu lunar rover separated from lander and started its march to explore the surface and the subsurface around the landing site. Yutu payload included an active particle induced X-ray spectrometer, a visible to nearinfrared (450-945 nm) imaging spectrometer and short-wave infrared (900-2395 nm) spectrometer (VNIS), and the LPR, together with a stereo camera and a navigating camera. The main goal of the radar was the estimation of the thickness of lunar regolith and the detection of the lunar subsurface structure along the traverse path. LPR is an ultrabroadband radar operating in time domain and equipped with two sets of antennas:

- 1) a pair of low-frequency monopole antennas (1150 mm long, 12-mm diameter, and separated by about 80 cm) with 60-MHz center frequency and 40-MHz bandwidth (i.e., 40–80 MHz), suspended 60 cm above the ground on the back of the rover (see Fig. 1);
- 2) a set of one transmitter and two receiver bowtie antennas (336 mm long, and 120 mm wide) operating at 500-MHz center frequency with 500-MHz bandwidth (i.e., 250–750 MHz).

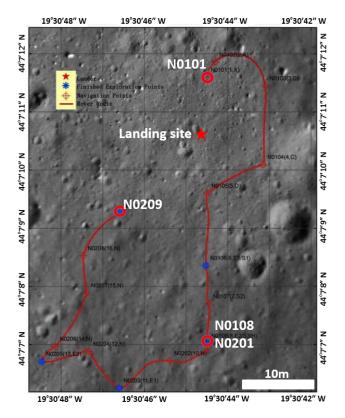


Fig. 2. Traverse path of the LPR on the moon. The red star represents the CE-3 landing site and the red curve means the rover route. The background topographic photograph was taken by the descending camera attached to the lander as it descended

These antennas were located at the bottom of lunar rover, about 30 cm above the ground, and separated 160 mm from each other [8]. We refer to these sets as channel 1 and channel 2, respectively. Channel 1 has a system gain of 152 dB and channel 2 of 133 dB. Detailed description of the radar system can be found in [8].

The LPR operational phases on the moon can be divided into three stages (see Fig. 2). In the first stage, the rover traveled 56 m in ten days, moving from navigation point N0101 to point N0108. During this phase, several radar parameters (e.g., system gain, time window, and attenuation settings) were tested to determine the performance of the system for the subsequent data acquisition. The second stage started at point N0201 and after 58-m reached point N0209; in this phase, the radar acquired data from both channels using the acquisition parameters previously set. Unfortunately, after this phase, the rover got stack in the moon soil and could not move any further. As a consequence in the last phase, the radar acquired data only at a fixed position (i.e., at point N0209).

# III. LPR CALIBRATION ON EARTH AND PREVIOUS STUDIES

LPR system was calibrated and tested on earth before launching, using an LPR prototype with a 1:1 simplified model of the CE-3 lunar rover, equipped with a GPS to track the radar profile location. Such prototype reproduces quite well performance of the system even though, in the actual

radar, specific isolators and filters to avoid the influence of communication signals between the lunar rover and the lunar lander were also installed [29]. The testing campaign was carried out by the Ground Research and Application System, NAOC, and Institute of the Electronics, Chinese Academy of Sciences, on three different types of ground: a glacier, a loess deposit, and an artificial lunar soil [29]. In the first two sites, both antennas (60 and 500 MHz) were tested, whereas in the artificial lunar soil, being as shallow as 15-70 cm, only the 500 MHz was used. However, as in this paper, we are interested in the 60-MHz data, we would not describe or comment any further the data collected with the higher frequency antennas. The best performance in terms of maximum penetration depth of the LPR prototype was achieved on the glacier experiment, where assuming a permittivity of 3.2, geological structures as deep as 152 m were recognized in the radar cross sections [29]. However, the data collected with the 60-MHz antennas on glacier also show, at various time depth (i.e., 1100, 2500, 3700, and 5800 ns), weak continuous signals which appear on the radar cross sections as quasi-horizontal features of constant amplitude [see Fig. 3(a)]. In particular, the shallowest one (1100 ns) seems to cut some geological structures visible on the right side of the section whereas the others seem to only superimpose to the noise. The same radar features are presented at similar time depths in the sections collected on loess deposit [see Fig. 3(b)] even though, in this case, the deepest interfaces (3700 and 5800 ns) are difficult to be recognized as they are almost totally buried in the noise. Nevertheless, none of the papers published on LPR calibration and testing have mentioned or discussed the presence and origin of such features.

Conversely, a detailed discussion on the two deepest features (3700 and 5800 ns), which are also visible on the radar sections collected on the moon (see Fig. 4), is presented in two different papers [27], [28]. Indeed, both studies assume that such features are real reflectors generated by some geological discontinuity. In particular, Xiao *et al.* [27] interpreted such reflectors as two layers representing different episodes of lava eruptions, whereas Zhang *et al.* [28] interpreted the reflectors as buried regolith layers between different basalt units. Note that in the data collected on the moon the two shallow radar features (at 1100 and 2500 ns) are not visible as totally buried in the strong ringing presented in the first 3000 ns of the time window (see [28, Fig. 5]).

A careful comparison between earth and moon data like, e.g., that illustrated in Figs. 3 and 4, seems to unambiguously indicate that such features are not real reflectors. In fact, it is highly improbable to find the same features in such different survey environments and operating conditions. Nevertheless, a simple qualitative analysis cannot be considered robust and definite evidence as, accounting also for the measurement uncertainties, some fortuitous coincidence in time depth cannot be excluded. Therefore, in the following, we will perform a detailed quantitative analysis to prove that such features do not correspond to real geological lunar structures but rather they are generated by the electronics and/or by the electromagnetic interaction between the system and the rover.

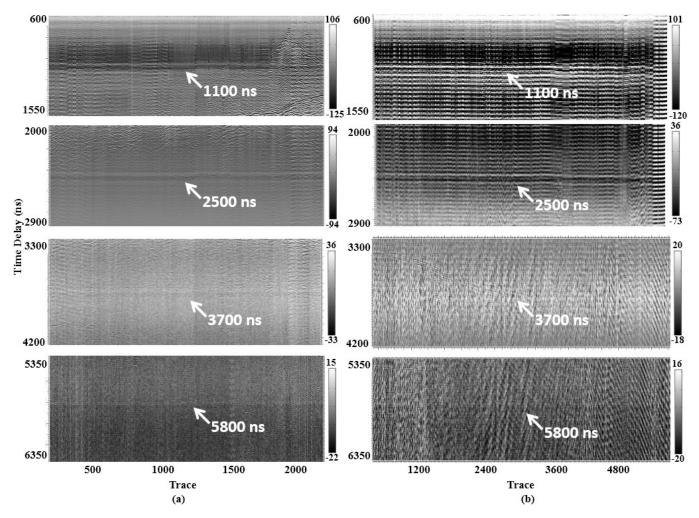


Fig. 3. Deep radar features detected on (a) glacier and (b) loess deposit. Note that the reflectors are located at the same time depth (indicated by the white arrow) but they are better visible on the glacier data as noise and ringing are lower, especially at larger depth, with respect to loess deposit data.

### IV. RADAR EQUATION AND SIGNAL PENETRATION

As the first step, we have tested the maximum signal penetration depth in lunar soil using the radar range equation and the parameters of the LPR at 60 MHz. Given the equation [31]

$$P_r = P_T \frac{G_{\text{Tx}} G_{\text{Rx}} \lambda^2 \sigma}{(4\pi)^3 D^4} e^{\left(-\frac{2\omega \tan \delta D}{\nu}\right)}$$
(1)

where

 $P_T$  transmitting power;

 $P_r$  receiving power;

 $G_{\text{Tx}}$  transmitting antenna gain;

 $G_{\rm Rx}$  receiving antenna gain;

 $\lambda$  wavelength in the medium;

velocity in the medium;

D depth of layer,  $D = v \times t/2$ ; t is the two way travel time.

 $\tan \delta$  loss tangent;

 $\sigma$  radar cross section;

 $\omega = 2\pi f;$ 

f = 60 MHz.

We can rearrange the terms to estimate system gain  $G_{\text{sys}}$  as

$$G_{\text{sys}} = \frac{P_T G_{\text{Rx}} G_{\text{Tx}}}{P_{\text{min}}} = \frac{(4\pi)^3 D^4}{\lambda^2 \sigma} e^{\frac{2\omega \tan \delta D}{D}}$$
(2)

where  $P_{\min}$  is the minimum detectable power at the receiver. The parameters in (2) are taken from Table I in [29] and the depth D is computed on the basis of the velocity v = 0.11 m/ns which is an average value for Mare Imbrium according to [6]. Furthermore, we have assumed  $\sigma = \pi D^2$ , which represents the maximum backscattering (total reflection) due to a flat and smooth interface. The choice of the value of  $\tan\delta$  is quite critical because it is not well constrained; therefore, we can only consider a wide range of values, i.e.,  $10^{-3}$ to  $8 \times 10^{-3}$ , as suggested by Ono et al. [6]. Fig. 5 illustrates the theoretical trend of the system gain values versus time depth for the above mentioned loss tangent range boundaries. The blue curve refers to the lower loss tangent ( $\tan \delta = 10^{-3}$ ) and the red curve refers to the highest value  $(8 \times 10^{-3})$ . It is evident that in this interval of values the dynamic range of channel 1 (152 dB) is suitable to detect bright reflectors up to a depth of about 5000 ns.

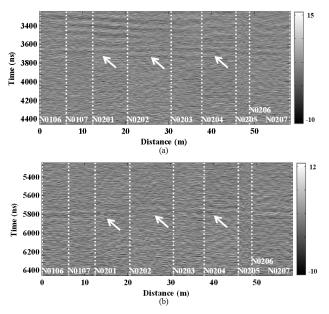


Fig. 4. Radar features on moon data collected at 60 MHz from navigation point N0106 to point N0207. (a) Radar features between 3300 ns and 4450 ns. (b) Radar features between 5300 ns and 6450 ns. The vertical dashed lines indicate different navigation points, which also refer to different days of acquisition. White arrows indicate two deep hypothetical reflectors.

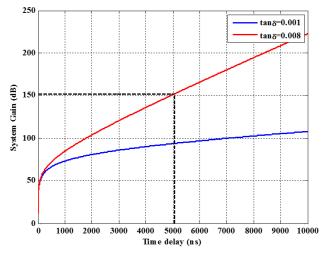


Fig. 5. Relationship between system gain and time delay computed by using (2). Blue curve indicates the most favorable case for signal penetration (low loss), and red curve indicates the worst case.

## V. DATA ANALYSIS AND RESULTS

#### A. Effect of Noise on Reflectors Detectability

As the second step, we have evaluated the reflector detectability on the moon data estimating the noise level and comparing such level with the signal amplitude of the deep reflectors. To this aim, we analyzed two data sets: the first one collected with the transmitter switched OFF (high-voltage-OFF mode) at navigation point N0209 (from 2014-02-15T22:51:38 to 2014-02-15T23:56:34), and the other with the transmitter ON (high-voltage-ON mode) acquired from navigation point N0101 to N0208. The estimation was made considering a subset of 150 traces extracted at N0201 (high-voltage ON) and N0209 (high-voltage OFF), respectively, and computing the noise level in terms of standard deviation

TABLE II

CALCULATED NOISE AT NAVIGATION POINTS N0201 (HIGH-VOLTAGE ON)

AND N0209 (HIGH-VOLTAGE OFF)

Model	σ(a.u.)
High-voltage-on	2.17±0.16
High-voltage-off	$2.16 \pm 0.16$

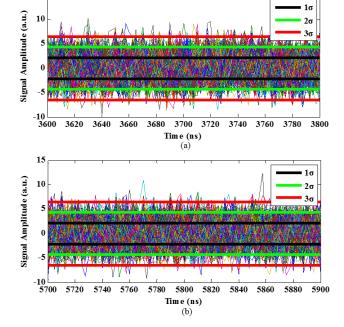


Fig. 6. Comparison between signal amplitude and noise level for the moon data. (a) Hypothetical reflector between 3600 and 3800 ns. (b) Hypothetical reflector between 5700 and 5900 ns.

of the signal amplitude. In high-voltage-OFF mode, the noise was calculated on the total time window (0–10000 ns) whereas in the high-voltage-ON mode only the bottom part of the time window (i.e., 6000–10000 ns) was considered to minimize the effects of weak reflectors presented above such time depth. As expected, the values computed for the two modes are very similar (see Table II) so we have assumed that for the noise level a standard deviation  $\sigma=2.16$ . In Fig. 6, it is reported the comparison between the amplitude values of all traces collected in the time windows 3600–3800 ns and 5700–5900 ns (i.e., the time intervals where the hypothetical reflectors have been detected) with the noise level computed at  $1\sigma$ ,  $2\sigma$ , and  $3\sigma$ . From Fig. 6, it is clear that the signal amplitude level in the two time intervals is quite similar and it barely exceeds the  $3\sigma$  noise level.

#### B. Time-Frequency Analysis Applying S-Transform

Further information about the nature of the radar features under investigation can be searched looking at their spectral content using different time–frequency representations (e.g., short time Fourier transform, Wigner–Wille distribution, or S-transform) [32]. In this paper, we have applied the S-transform as it represents a good compromise between

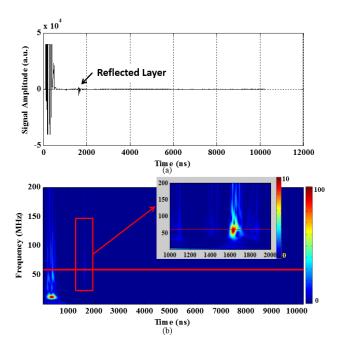


Fig. 7. Example of the S-transform based on glacier data. (a) One trace of radar data. (b) S-transform of (a).

frequency–time resolution and simple spectrum interpretation. This technique is conceptually similar to the continuous wavelet transform and it is based on a moving and scalable localizing Gaussian window. S-transform S(t, f) of signal x(t) is defined as

$$S(t,f) = \int_{-\infty}^{+\infty} x(\tau) \frac{|f|}{\sqrt{2\pi}} e^{\frac{-f^2(t-\tau)^2}{2}} e^{-i2\pi f \tau} d\tau$$
 (3)

where f is the frequency.

First of all, we tested such technique on the data collected on the glacier, analyzing the signals reflected by natural discontinuities (i.e., the dipping reflector detected by Zhang et al. [29]) presented in the ice layer at a time depth between 1400 and 1800 ns. Fig. 7 shows the results of the S-transform analysis applied to the entire time window of a single trace. More specifically, Fig. 7(a) illustrates the time trace and Fig. 7(b) illustrates the time-frequency trace. In the latter, the insert box shows the frequency content of the reflector at about 1600 ns, which exhibits a spectrum centered at 60 MHz in full agreement with the central frequency of channel 1 transmitting antenna. We then applied the same procedure to the complete glacier data set [Fig. 8(a)] and we extracted from the S-transform only the 40-80MHz components in order to generate the time-amplitude image reported in Fig. 8(b). The comparison between these images shows that the natural discontinuity inside the glacier produces echoes that preserve the original frequency content emitted by the antenna.

As the last test on the glacier data, we have selected the time interval around the four hypothetical reflectors (see Section III) to study the frequency content of such features. However, because the amplitude of these features is of the same order of magnitude of the noise level (see Fig. 9), as also noticed for the moon data (see Section III), the S-transform has

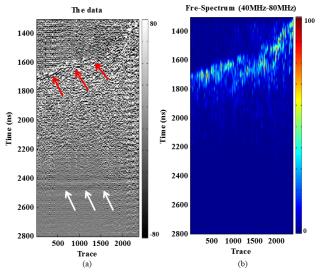


Fig. 8. S-transform of the glacier data. (a) Chosen part. (b) S-transform result with frequency spectrum (40–80 MHz).

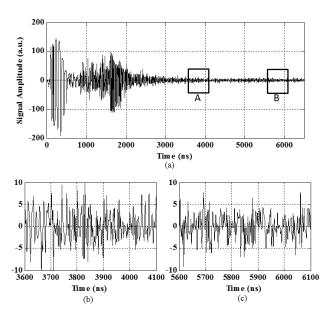


Fig. 9. Example of radar trace collected on the glacier. (a) Boxes labeled by A and B indicate the position of the deepest events also present on the moon data. (b) Enlarged view of data in box A. (c) Enlarged view of data in box B.

been applied on a single stacked trace averaging 1000 traces. Fig. 9 illustrates, as an example, the level of the signal on a radar trace in the time interval around the position of the two deepest hypothetical reflectors, whereas Fig. 10 shows the stacked trace (a), the signal in the time intervals of interest (b), and the corresponding S-transform (c). Note that in each graph of Fig. 10(b), the red curve is a sinusoidal signal having time position and frequency given by the coordinates (f, t) of the maximum value of the S-transform illustrated in Fig. 10(c). The spectrum central frequency of all four hypothetical reflectors is substantially the same (about 12 MHz), and differently from the previous case (see Figs. 7 and 8); it is well below the bandwidth of the transmitted signal (40–80 MHz).

Finally, the same procedure described above was applied to the moon data, computing the average trace from 1000 traces collected at navigation point N0201. Fig. 11 illustrates the

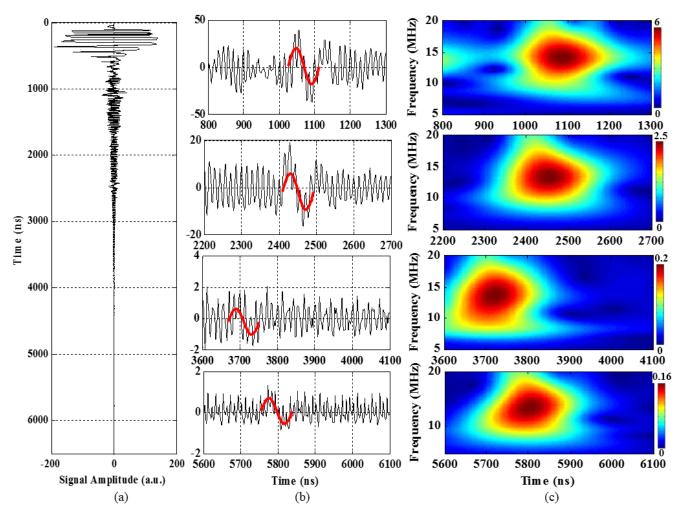


Fig. 10. Results of the S-transform applied to the averaged trace computed from glacier data. (a) Position of the time intervals where the hypothetical reflectors are located. (b) Enlarged view of the time intervals. (c) Relevant S-transform images.

results of the analysis in a similar fashion to those in Fig. 10 but only for the two deepest hypothetical reflectors as the shallow ones cannot be extracted from the ringing (see Section III). The results are in good agreement with those obtained for the glacier as the time depth and the spectrum central frequency (about 12 MHz) values of the features are quite similar.

#### VI. DISCUSSION

As highlighted in Section III (see Figs. 3 and 4), shape, amplitude, and position of the hypothetical reflectors appear to indicate that they are stationary disturbances always present in the radar cross sections and well visible only if some specific processing is applied to enhance them above the noise level (see supporting information of [28]). In particular, a careful comparison between Figs. 3 and 4 highlights that the shape of such features on the moon data is slightly curved if compared with those present, at the same time depth, in the glacier and loess data. This fact is probably due to the difference in terms of acquisition time of the radar data set; indeed, the data collected on the glacier and loess terrains were made in a single day, whereas the data collected on the moon refer to 15 different days and thus different

operating modes of the rover and of the other instruments. As a consequence, the lunar radar cross section is an assemblage of different data sets that do not properly align, even though inside single blocks of data (i.e., between two subsequent navigation points) the radar features are essentially horizontal (see also [28, Fig. 5]).

The theoretical computation performed assuming an ideal reflector and considering the expected values for the loss tangent in Mare Imbrium has shown that it cannot be excluded that the LPR dynamic range (152 dB) could be sufficient to detect reflectors as deep as 300 m. However, the quantitative analysis performed on the signal amplitude supports the hypothesis that the deep radar features under question do not come from real geological structures. The analysis performed on the noise highlights that there is no significant difference in terms of amplitude level between the noise and the signals in the time windows associated with the two hypothetical reflectors. In fact, the signal amplitude level for the lunar deep reflectors is of the order of  $3\sigma$ ; however, similar values can also be found for the data collected on the glacier. Furthermore, the similarity between the amplitude levels of the signals coming from different depths poses serious questions

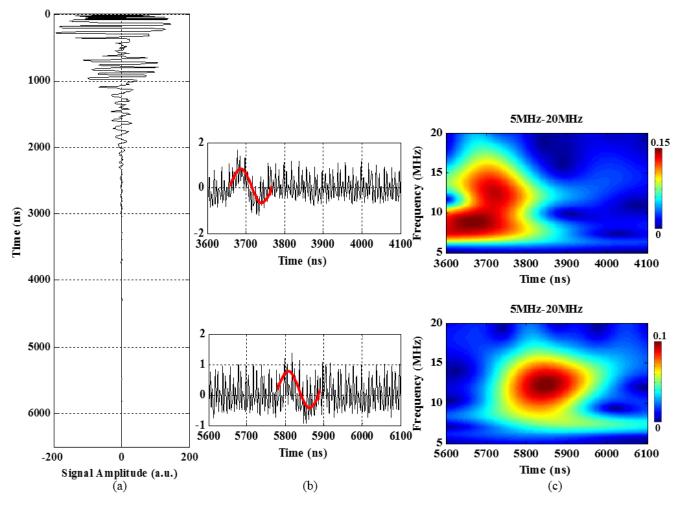


Fig. 11. Results of the S-transform applied to the averaged trace computed from moon data. (a) Position of the time intervals where the hypothetical reflectors are located. (b) Enlarged view of the time intervals. (c) Relevant S-transform images.

about the reliability of the data. Indeed, if we assume that the signals coming from the first deep reflector (at about 3700 ns) are real and their small amplitude is due to the propagation in the soil (nonnegligible attenuation), it is physically unlikely that the second reflector (at about 5800 ns) could maintain the same amplitude after over 2000 ns of propagation.

The S-transform analysis adds another important piece of information to clarify the nature of radar features interpreted by Xiao et al. [27] and Zhang et al. [28] as real reflectors. The main outcome of such analysis is the presence of a systematic frequency shift in the radar signals regardless survey location (glacier or lunar terrain) and time delay. Indeed, for both terrestrial and lunar radar data, the values of the central frequency (10-15 MHz) of the signals hypothetically coming from deep interfaces are well below the central frequency of the transmitted signals (60 MHz). On earth, it is well known that a radar signal can experience an apparent shift in frequency when it propagates through a soil acting as low-pass filter [33], [34]; such shift is ruled by the relationship between attenuation and frequency. In the case of signals propagating through the ice (glacier survey), in which the attenuation is frequency independent, no frequency shift should be expected

and any signal coming from a real reflector should preserve the original frequency content. The analysis of signals coming from the interface located between 1400 and 1800 ns in the glacier radar cross section is in full agreement with this statement (see Fig. 7). On the contrary, the signal coming from the hypothetical deep reflectors has a frequency content which cannot be justified by the propagation in the subsurface. Differently from ice, the materials composing Mare Imbrium terrains can act as low-pass filter and can produce a shift in frequency as, according to our knowledge about lunar materials, the radar attenuation is frequency dependent [35], [36]. However, such shift can be appreciated only if the attenuation (or loss tangent) in the material and the time depth of the reflector are large enough. A way to verify if the shift could be real is to compute the value of the loss tangent  $(\tan \delta)$ that would have produced such effect. In practice, considering the two hypothetical reflectors (at about 3700 and 5800 ns),  $\tan\delta$  can be estimated applying (11) in [36] and assuming a bandwidth of 40 MHz. For the shallower reflector, we found  $\tan \delta_{3700 \text{ ns}} = 1.6 \times 10^{-2}$  and for the deepest one  $\tan \delta_{5800 \text{ ns}} =$  $0.8 \times 10^{-2}$ ; both values are larger than those expected for Mare Imbrium terrains (see Fig. 5). This result poses again a

question about the reliability of the radar data. In fact, if we assume that the computed loss values were real, the maximum penetration depth achievable by the system on the moon would be lower than 3700 ns (see Section IV) and no deep signal could be actually detected.

Finally, despite a total disagreement about the interpretation of the radar features between this paper and [27], our spectral analysis results are strongly supported by Zhang *et al.* [28]. In fact, Zhang *et al.* [28] found that the only way to extract the noise from the signals associated with the hypothetical deep reflectors is to apply a bandpass filter between 4 and 30 MHz, i.e., outside the frequency band of the transmitted signal.

#### VII. CONCLUSION

In this paper, we have analyzed the characteristics of the radar signals collected on the moon by LPR using channel 1 antennas. The scope of the work was to clarify the origin of specific radar features visible in the radar cross sections and previously interpreted as real reflectors associated with the layering structure of the lunar subsoil. Such analysis was based on three different approaches: 1) a qualitative comparison between the radar data collected on the earth and on the moon using channel 1 LPR system; 2) the amplitude of the radar signal versus the background noise level; and 3) the comparison between the signal frequency content of the transmitted and hypothetically reflected signals. We found that the overall results are robust, fully consistent, and totally unambiguous. Therefore, we can conclude that the deep radar features are not real reflectors rather they are signal artifacts superimposed to the radar traces. Regarding the origin of these artifacts at the moment, we can only speculate. Our study has shown that these signals are presented, with almost identical characteristics, in both terrestrial and lunar data, suggesting that they are probably generated by the electronic of the system and/or the radar-rover coupling. In 2018, a new opportunity to study the moon subsurface will be offered by CE-4 mission, which will be equipped with the same radar as CE-3. It would be of paramount importance to fully understand the nature of such strong disturbance and, possibly, eliminate it to obtain a reliable view of the lunar shallow geology.

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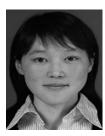
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